



Online Low Temperature Plasma Seminar
6th September, 2022

<https://theory.pppl.gov/news/seminars.php?scid=17&n=oltp-seminar-series>

Cathodic arc plasma science and applications: The known unknown and unknown known

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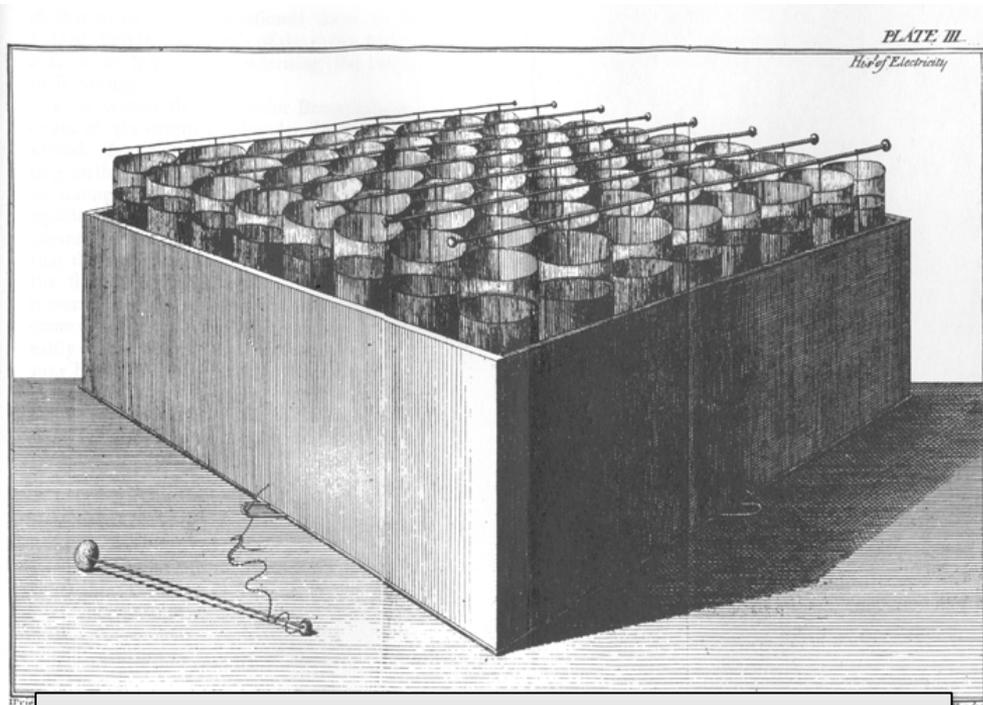
Colleagues, past and presents, in Berkeley, Leipzig and beyond, whose contributions are mentioned, are gratefully acknowledged!

Cathodic arc plasma science and applications: The known unknown and unknown known

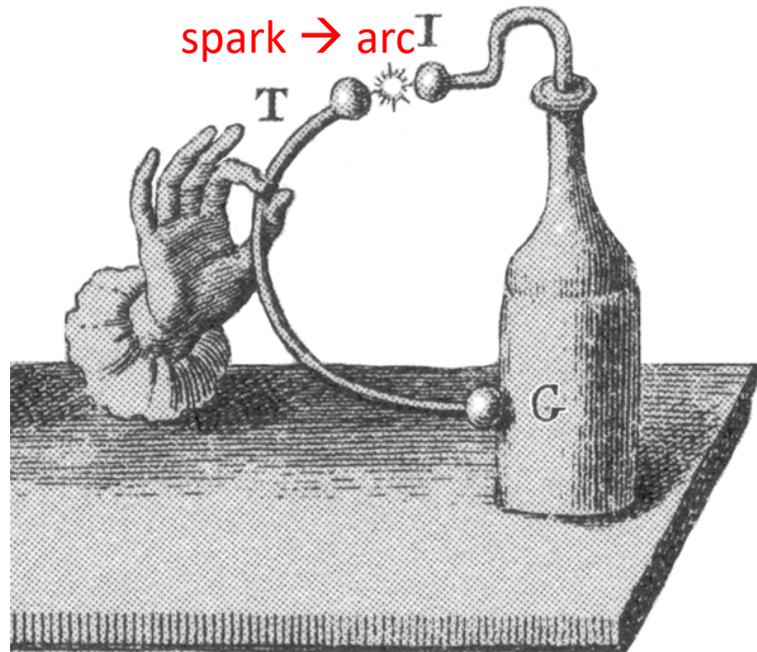


→ to be on solid ground, focus on arc plasma physics:
there are many (unknown) knowns and (known) unknowns ...

Pulsed Arcs in Air, (unintentional) Energetic Deposition



Battery of Leyden Jars (capacitors), ca. 1760



Alessandro Volta, ca. 1765

J. Priestley, *The History and Present State of Electricity*, London 1766

A. Anders, *IEEE Trans. Plasma Sci.* 31, 1052 (2003), and "Cathodic Arcs", Springer, NY 2008.

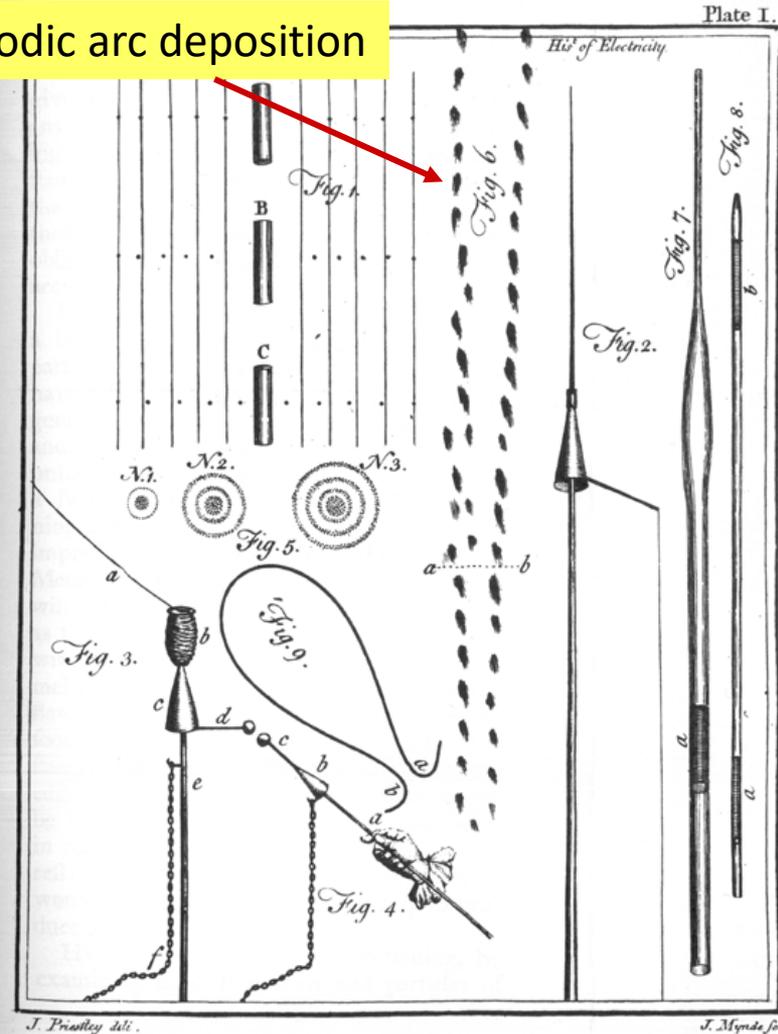
1766 - Joseph Priestley

- ❑ Cathodic arc coating (in air)
- ❑ discharge of a bank of Leyden jars through a brass chain
- ❑ arcs between each link of the chain
- ❑ deposit on glass is well adherent
- ❑ observed Newton's rings (interference, oxide films)
- ❑ found black coating (copper oxide)

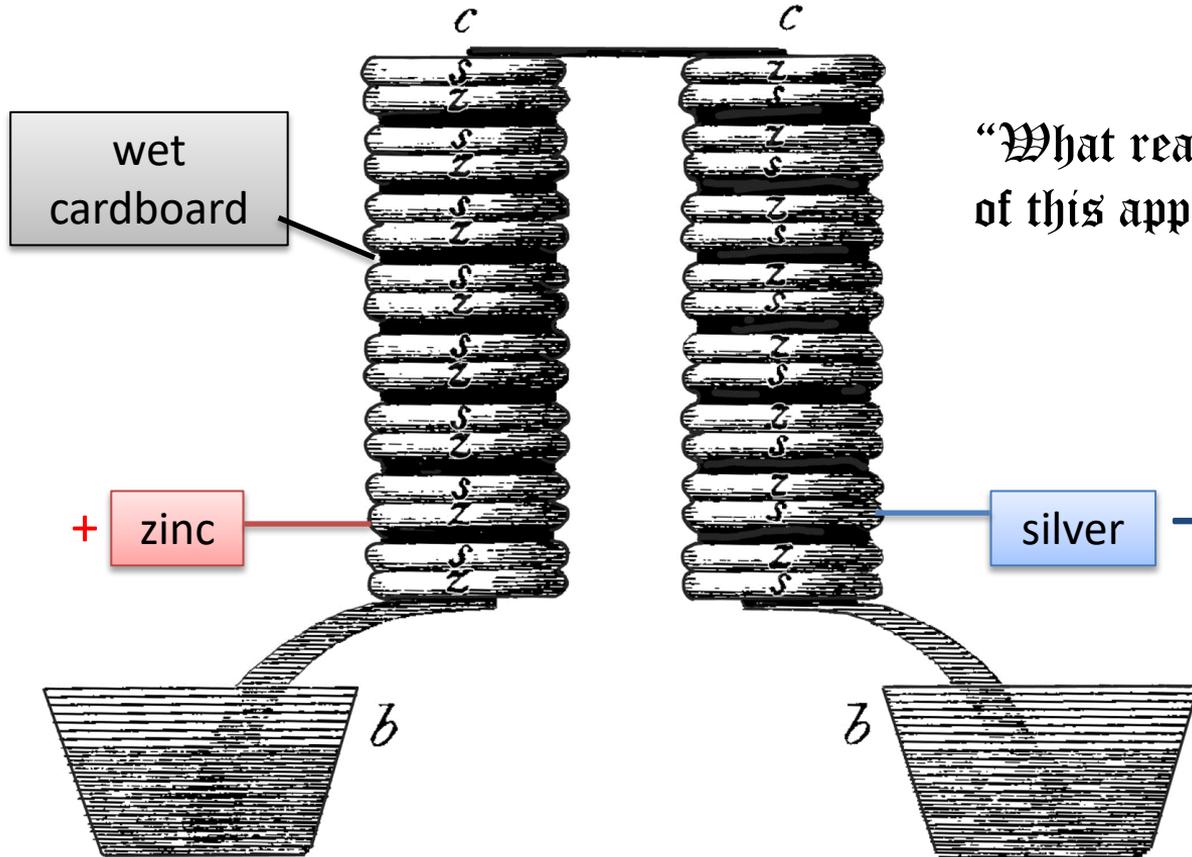
J. Priestley, *The History and Present State of Electricity*, London 1766

A. Anders, *IEEE Trans. Plasma Sci.* 31, 1052 (2003), and "Cathodic Arcs", Springer, NY 2008.

Cathodic arc deposition



Volta, 1800: The “Pile” – A New Electrical Energy Source



“What really increases the electric power of this apparatus is the number of plates”
Volta, 1800

A. Volta, Phil. Mag. VIII
(1800) 289.

1802/03 Vasilii Petrov

- ❑ First continuous arcs in air and at low pressure using an “enormous battery” of 4200 copper-zinc plates
- ❑ published only in Russian, his work was unknown or ignored

(on arc history) A. Anders, IEEE
Trans. Plasma Sci. 31 (2003) 1052

ИЗВѢСТІЕ

О

ГАЛЬВАНИ - ВОЛЬТОВСКИХЪ

О П Ы Т А Х Ъ,

которыя производилъ

Профессоръ Физики Василій Петровъ,

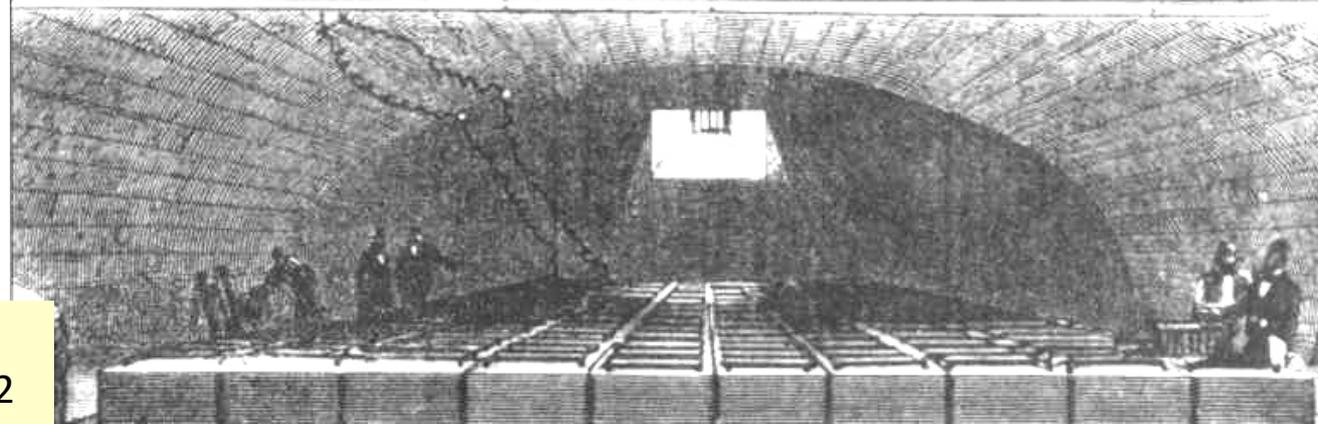
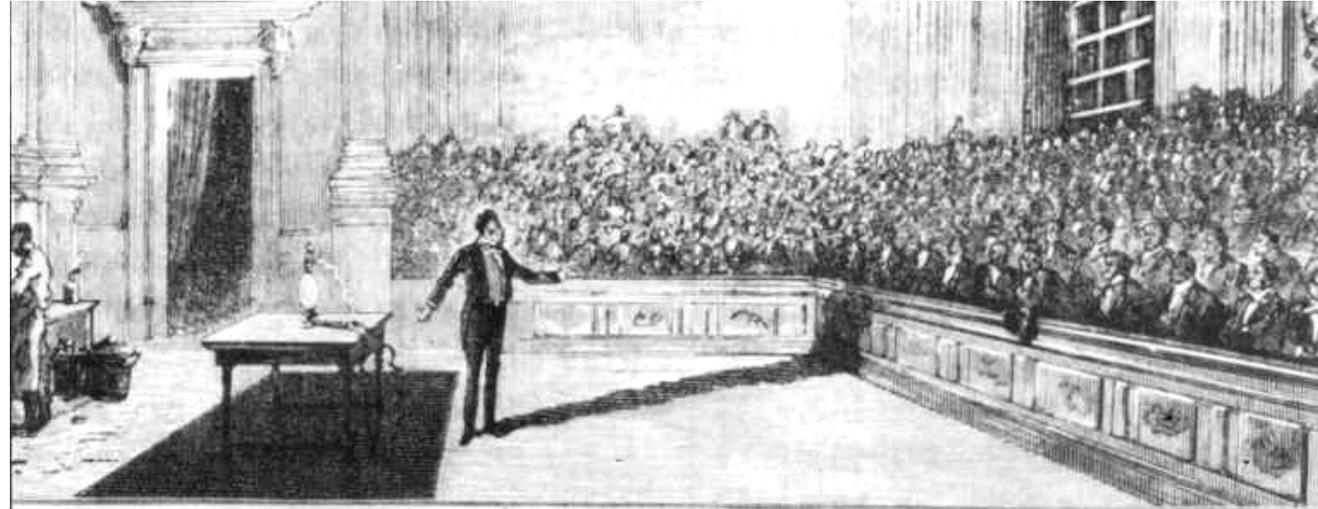
посредствомъ огромной наипаче бат-
терей , состоявшей *иногда* изъ 4200
мѣдныхъ и цинковыхъ кружковъ, и на-
ходящейся при Санкт - Петербургской
Медико - Хирургической Академіи.

ВЪ САНКТ-ПЕТЕРБУРГѢ ;

Въ Типографіи Государственной Ме-
дицинской Коллегіи , 1803 года.

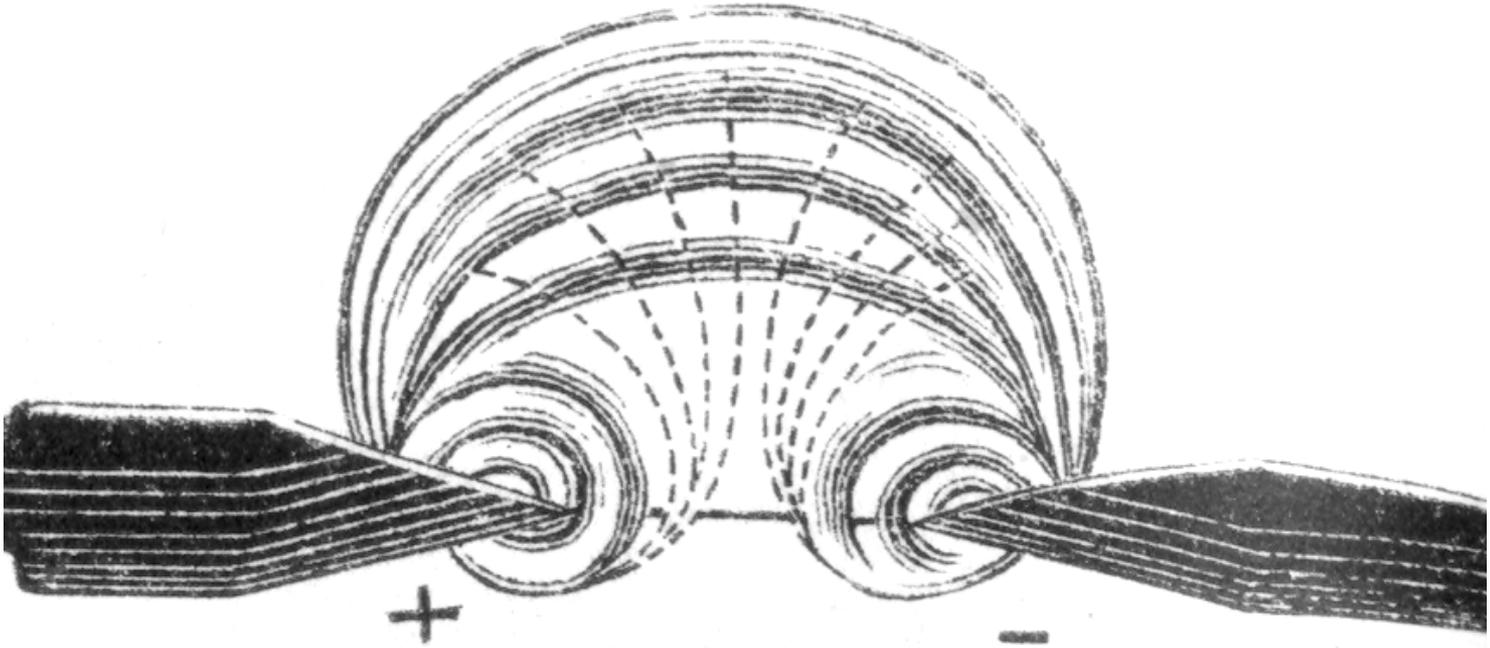
1802...1821 – Arc Discharge Demonstrations

- Humphry Davy at the Royal Institution London: Continuous arcs in air and in low-pressure vessels (1809)



(on arc history) A. Anders, *IEEE Trans. Plasma Sci.* 31 (2003) 1052

Arc Discharge



G. Wiedemann, Die Lehre von der Elektrizität, Braunschweig 1885

Important: Improved vacuum pumps of the 19th century

- Geissler's idea (1855) to use mercury as a „piston“ to pump air: then the base pressure is given by mercury's vapor pressure (0.5 Pa at room temp.)
- used by Edison to pump his light bulbs

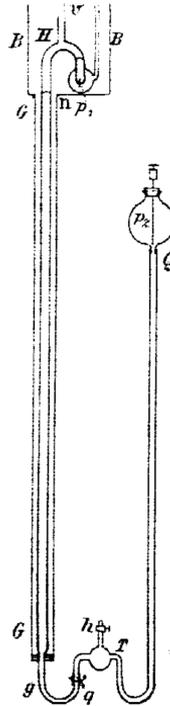


Heinrich Geissler (1814–1879)

Mercury Arc

Leo Arons (1892)

"Despite the continual impression of the appearance of light, the discharge is always discontinuous; the discontinuity is easily recognized by a telephone receiver, which is connected next to the apparatus in addition to a large capacitor. The discontinuity is due to the fact that, despite the measured low voltage at the electrodes, an electromotoric force almost three times as large is required to produce the arc; 50 volts are just sufficient."



Jahrg. 11.

Nr. 6.

Verhandlungen der Physikalischen Gesellschaft zu Berlin.

Sitzung vom 21. October 1892.

Vorsitzender: Hr. A. KUNDT.

Hr. W. Jäger berichtete auf Grund von gemeinsam mit
Hrn. D. Kreichgauer angestellten Versuchen

Ueber den Temperaturcoefficienten des Quecksilbers.

Hr. L. Arons sprach dann

Ueber einen Quecksilberlichtbogen.

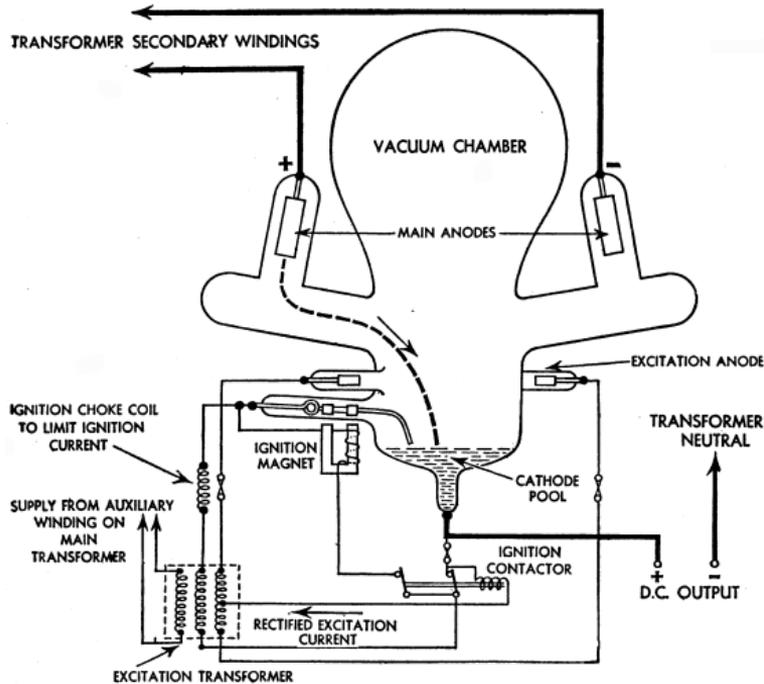
Mit Versuchen über die Gasentladung beschäftigt, fand ich eine ausserordentlich einfache Methode, einen intensiv leuchtenden, lang andauernden Lichtbogen zwischen Quecksilberelectroden herzustellen, welcher nur verhältnissmässig geringe electromotorische Kräfte zu seiner Unterhaltung erfordert und keine der lästigen Eigenschaften von Lichtbögen zwischen Metallelectroden besitzt.

L. Arons, Verhandlungen der Physikalischen Gesellschaft zu Berlin, 11 (1892) 767-771.

L. Arons, (Wiedemann's) Annalen der Physik und Chemie, 58 (1896) 73-95.

Mercury arc rectifiers

also known as *Cooper-Hewitt* or *Hewittic rectifiers* were extensively used to provide DC in high power applications, powers ranging from kW to MW, with voltages from 110 V to 30 kV.



Arc rectifier used in a NY substation for 90 years

Cathode Spots



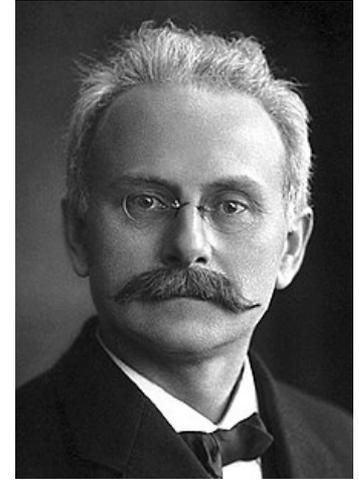
Fig. 5.



Fig. 6.



Fig. 7.



Johannes Stark (1874-1957)
(same person who found
broadening of spectral lines
due to electric fields)

Arc spot on mercury cathode , displacement when magnetic field is present

Cathode Spots:

Q1: Why do they form?

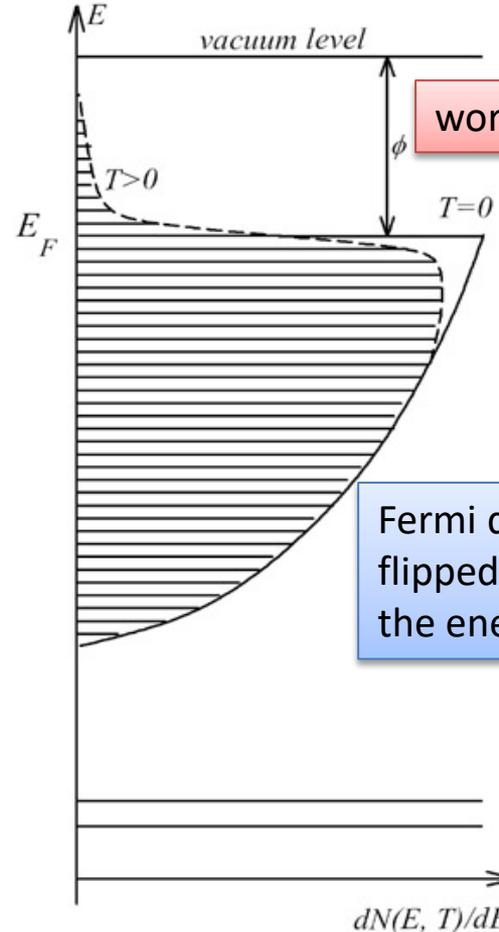
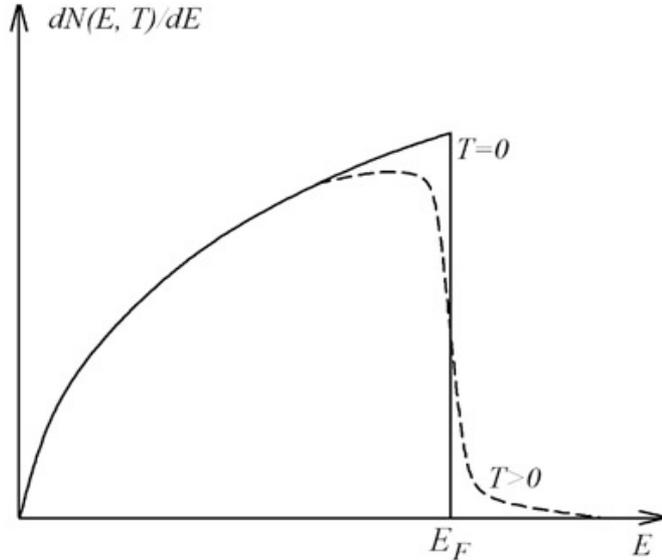
Q2: Why are they non-stationary?

A1: Because thermionic and field emission mechanisms are nonlinear with surface temperature and electric field

A2: Because they change the emission environment in nonlinear ways

Electrons in Metals: Fermi Distribution

Fermi-Dirac distribution function



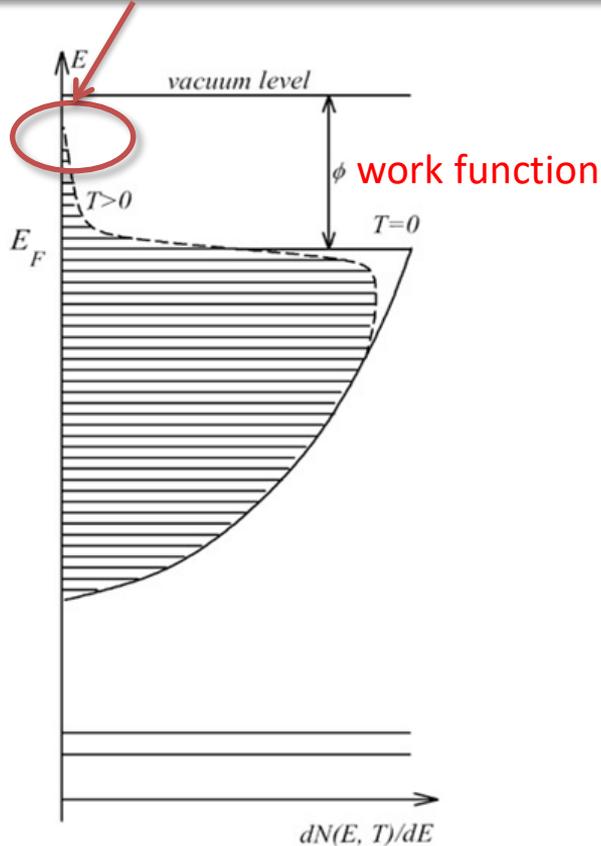
work function

Fermi distribution function flipped and rotated as to have the energy axis pointing up

$$\frac{dN(E, T)}{dE} = \frac{4\pi (2m_e)^{3/2}}{h^3} \frac{E^{1/2}}{1 + \exp\left(\frac{E - E_F}{kT}\right)}$$

Thermionic Emission from Hot Metal

those hot electrons, the tail of the distribution, can go OVER the potential barrier



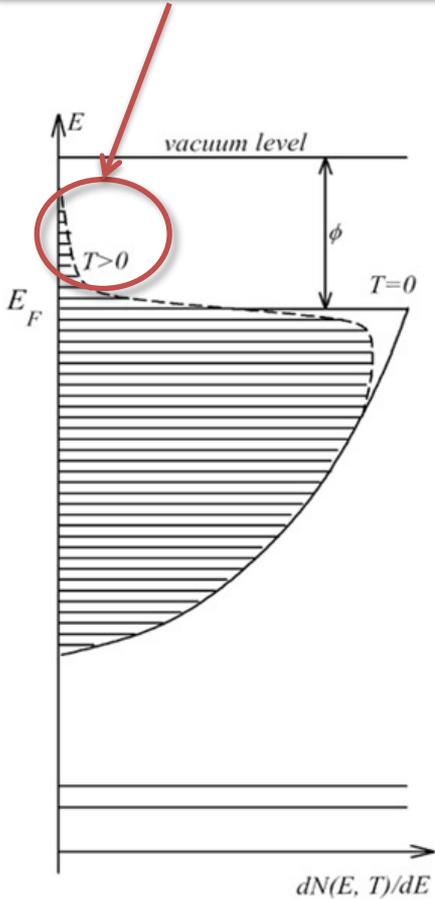
Richardson-Dushman equation

$$j_{thermionic} = AT^2 \exp\left(-\frac{\phi}{kT}\right)$$

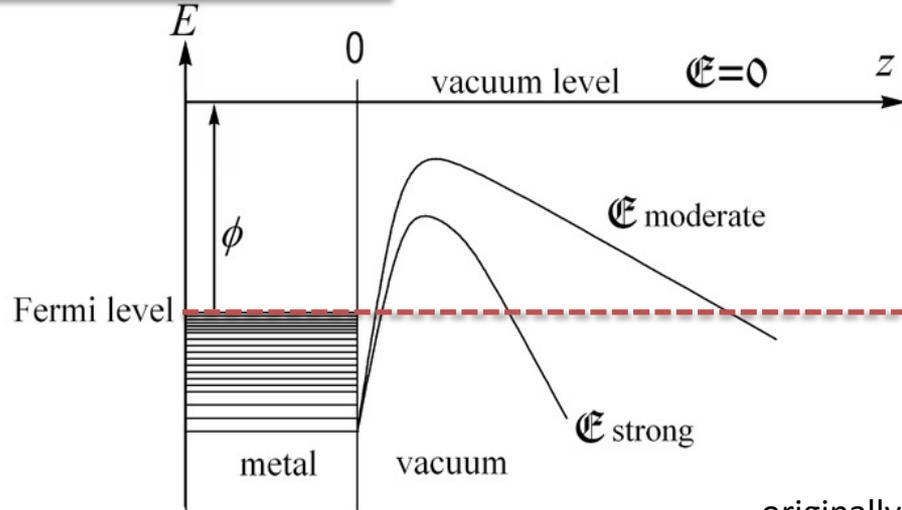
with the Richardson constant

$$A = \frac{4\pi em_e k^2}{h^3} = 1.202 \times 10^6 \text{ A/m}^2 \text{K}^2$$

less energetic electrons can tunnel THROUGH the potential barrier, if E-field is strong



Field Emission of Electrons



reduction of work function

$$\phi_S = \phi - \left(\frac{e^3 \mathcal{E}}{4\pi\epsilon_0} \right)^{1/2}$$

originally treated by Fowler-Nordheim (is rather complicated, even for $T = 0$)

$$j_{FN}(\mathcal{E}) = \frac{e^3}{8\pi h \phi t^2(y)} \mathcal{E}^2 \exp \left[-\frac{8\pi \sqrt{2m_e} \phi^3}{3eh\mathcal{E}} v(y) \right]$$

$$y \equiv \sqrt{\frac{e^3 \mathcal{E}}{4\pi\epsilon_0}} \frac{1}{\phi}$$

elliptical functions

Thermo-Field Emission of Electrons

Combining classical “**over** the barrier” with quantum-mechanical “**through** the barrier” gives a somewhat complicated expression for the emission current density:

$$j_{TF}(T, \mathcal{E}, \phi) = \frac{4\pi m_e kT}{h^3} \left\{ \int_{-W_A}^{W_I} \frac{\ln \left[1 + \exp \left(-\frac{E_z + \phi}{kT} \right) \right]}{1 + \exp \left[\frac{8\pi (2m_e)^{1/2} v(y) E_z^{3/2}}{3he} \right]} dE_z + \int_{W_I}^{\infty} \ln \left[1 + \exp \left(-\frac{E_z + \phi}{kT} \right) \right] dE_z \right\}$$

The rest of the definition of all symbols would take 2 more slides...
they can be found in chapter 3 of

Thermofield-Emission to Explosive Emission

Dynamics of cathode spots:
carbon arc as seen with a high-speed framing camera
arc current up to 5 kA, 200 ns exposure time per frame

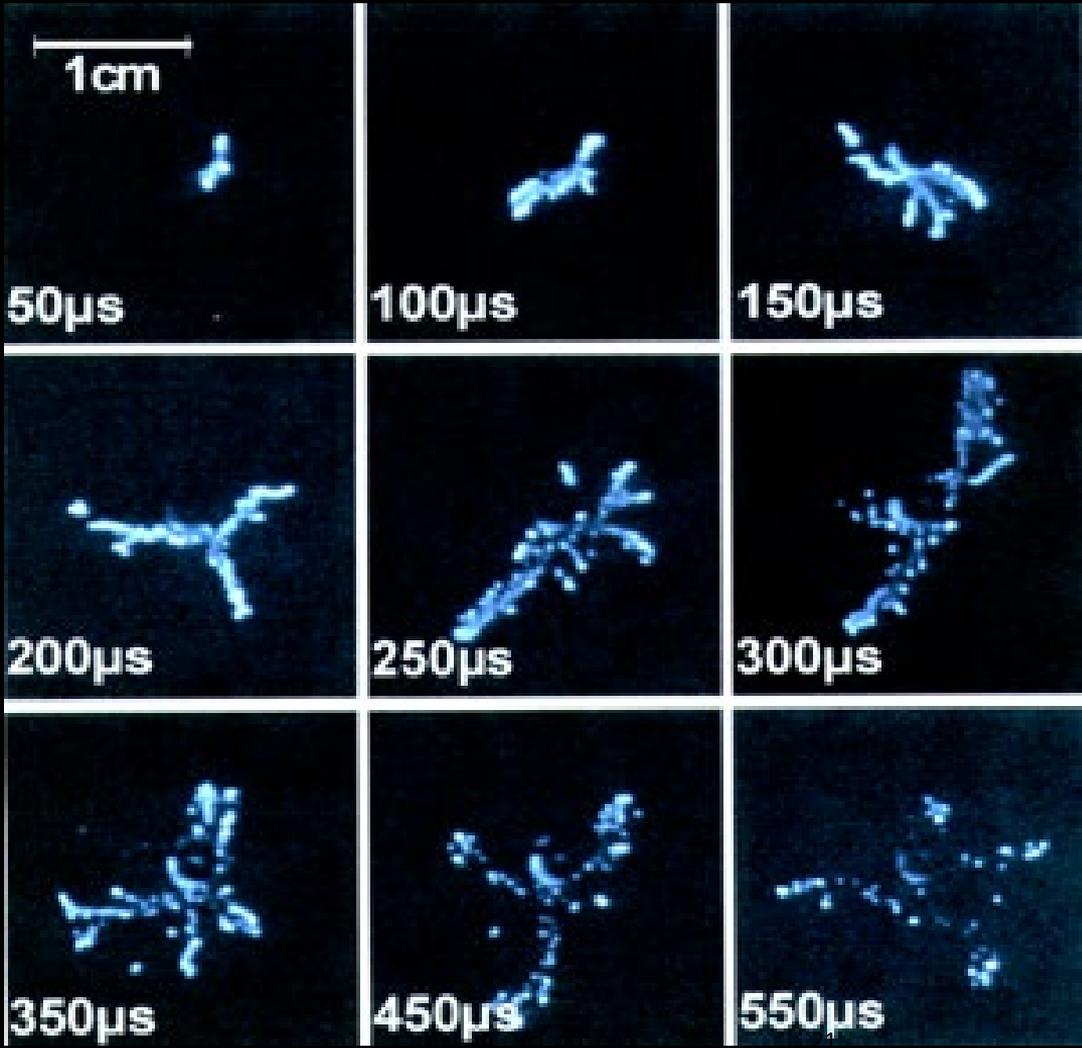
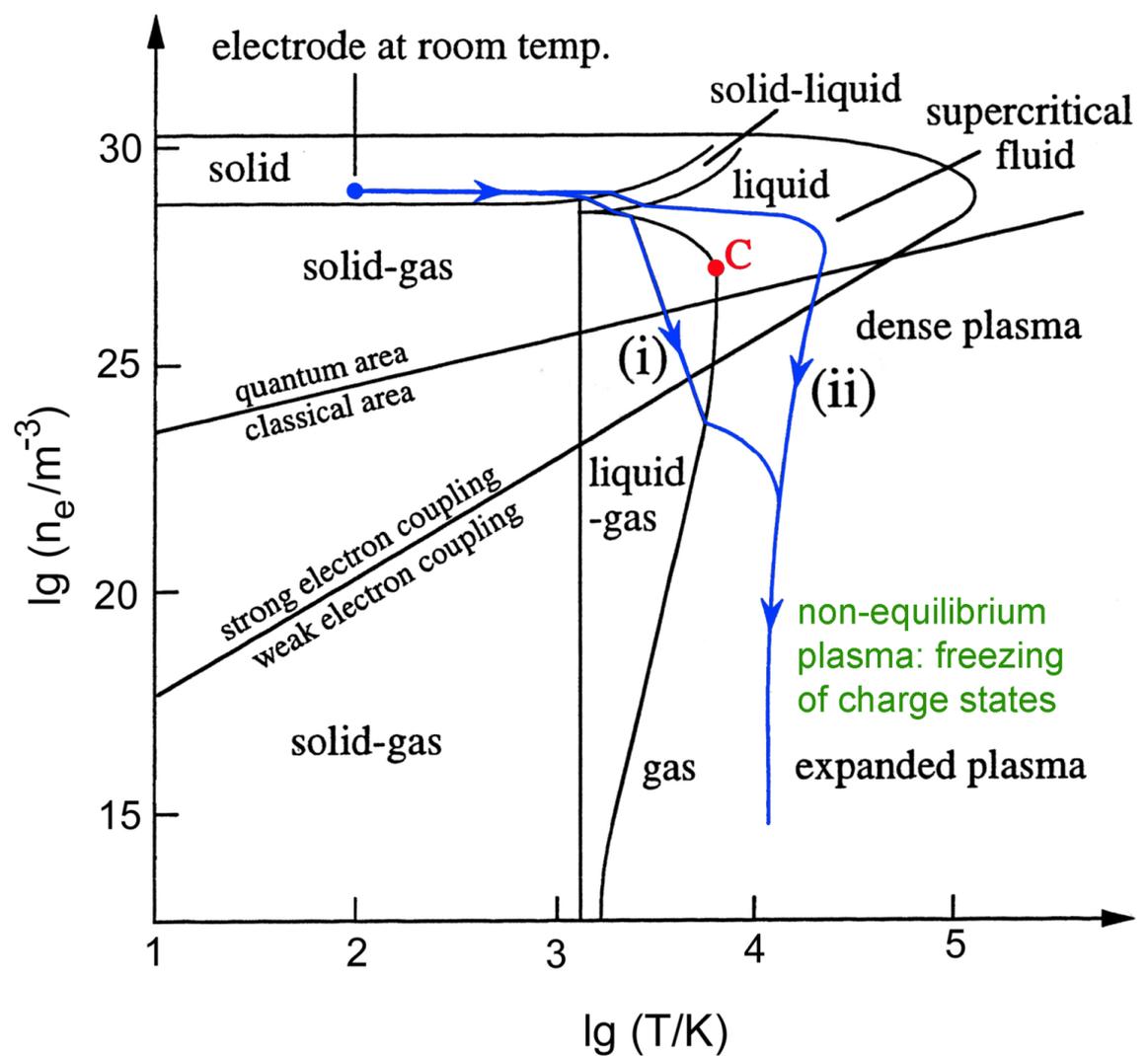
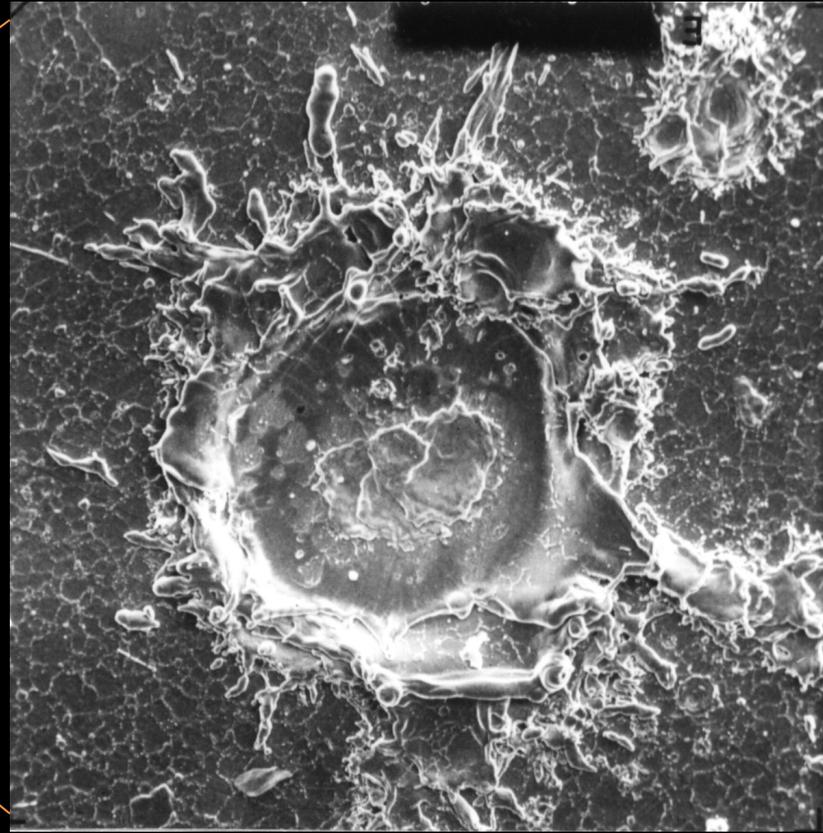
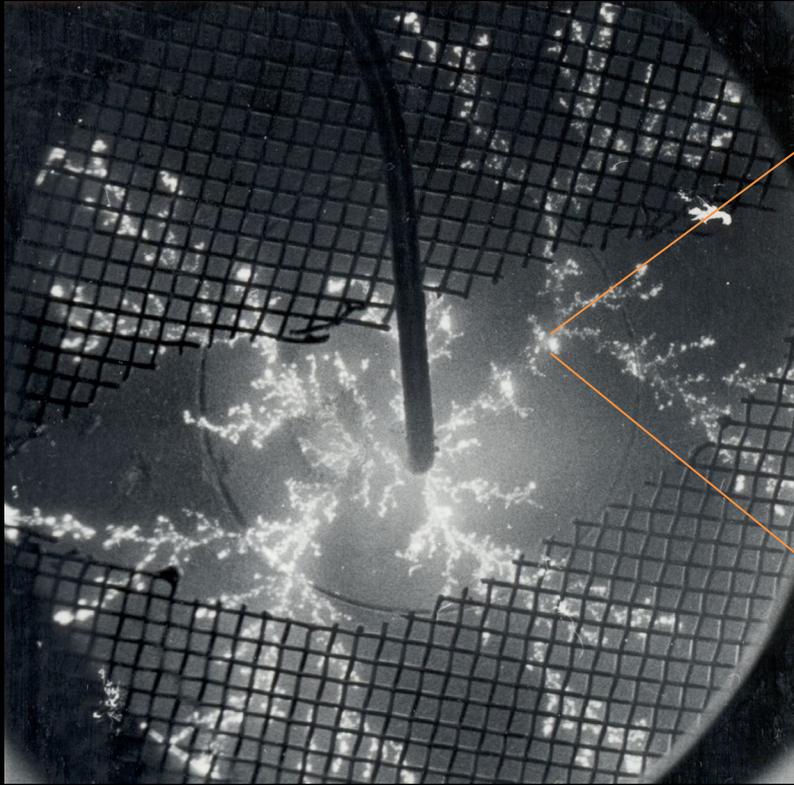


Figure: Courtesy of P. Siemroth and T. Schülke (about 2000).

Path of the Cathode Material in the n - T Phase Diagram

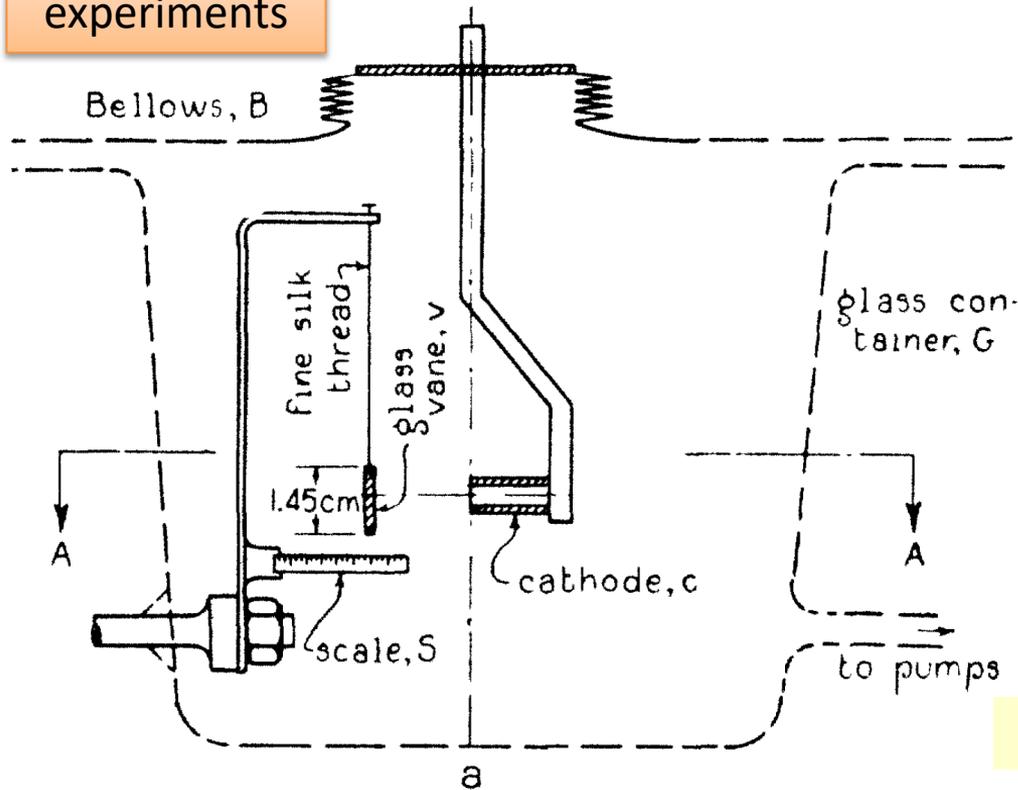


Effects of high power density in cathode spot



Supersonic ion flow

Pendulum experiments



→ Cu ion velocity = 1.6×10^4 m/s

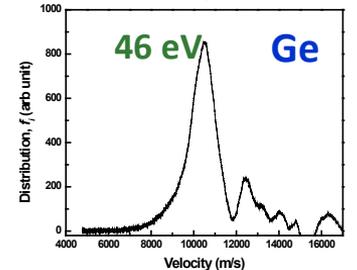
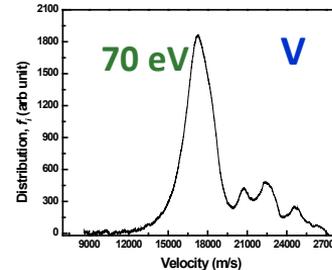
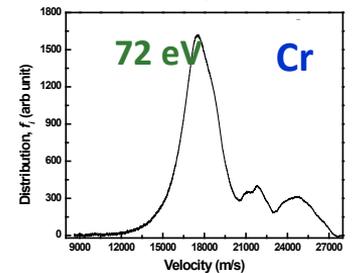
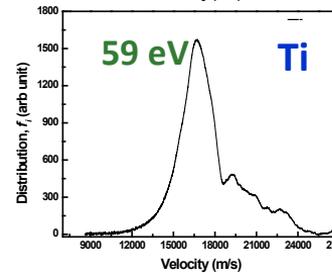
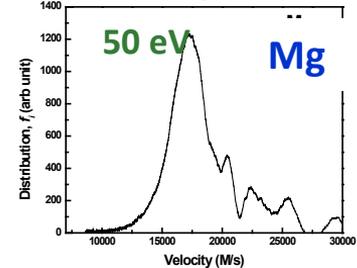
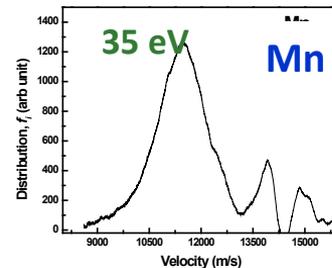
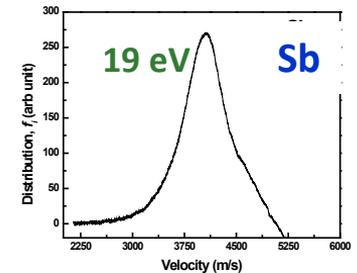
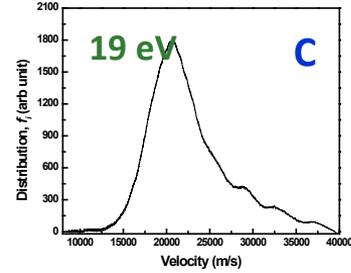
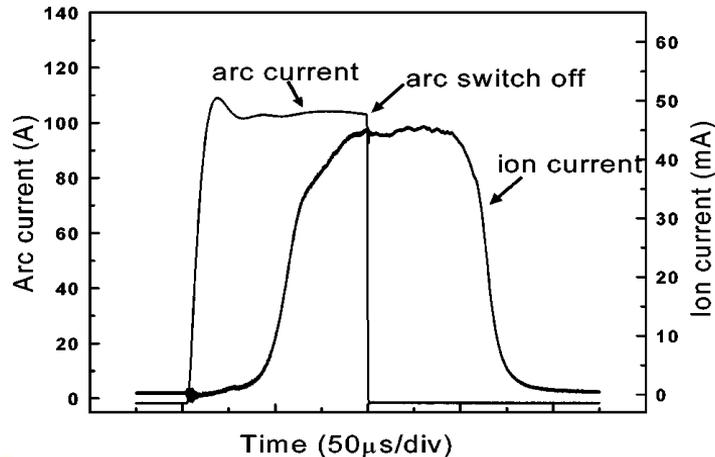
Is correct value but was criticized because he said "temperature" 500,000 K

R. Tanberg, Phys. Rev., 35 (1930) 1080-1089.

Both these methods gave a vapor velocity of the order of 16×10^5 cm/sec. A temperature of around $500,000^\circ$ K results when this value for the cathode vapor velocity is substituted for c in the equation: $\frac{1}{2}mc^2 = 3KT/2$.

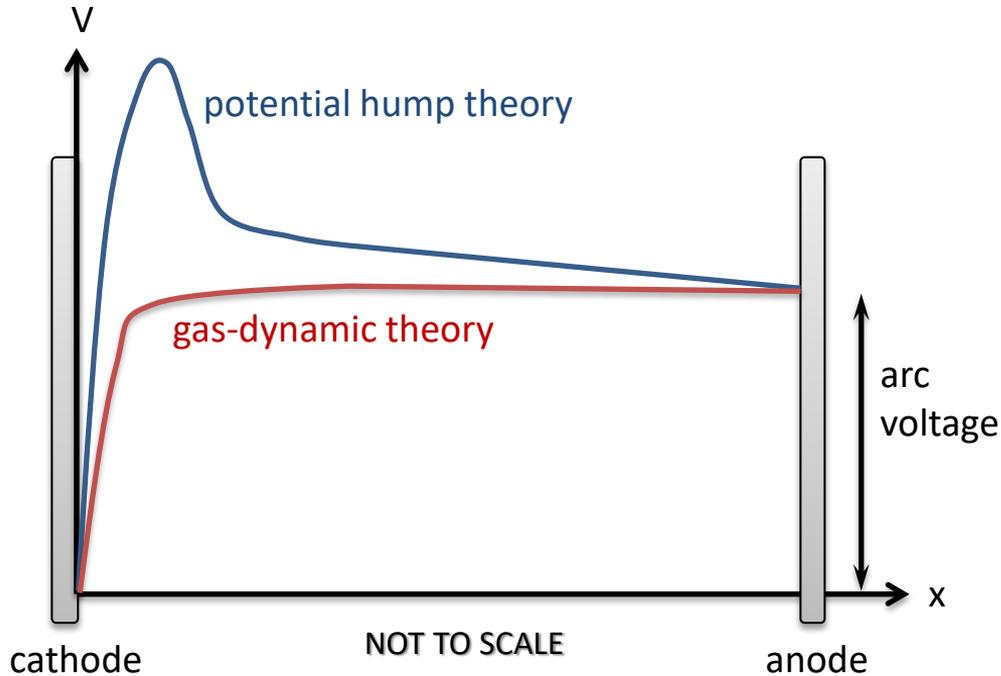
Energy Distribution Functions for Cathodic Vacuum Arcs

- Switch-off TOF method
- Distributions show **one** large peak
- → kinetic energy is almost independent of ion charge state; a pressure driven ion acceleration mechanism



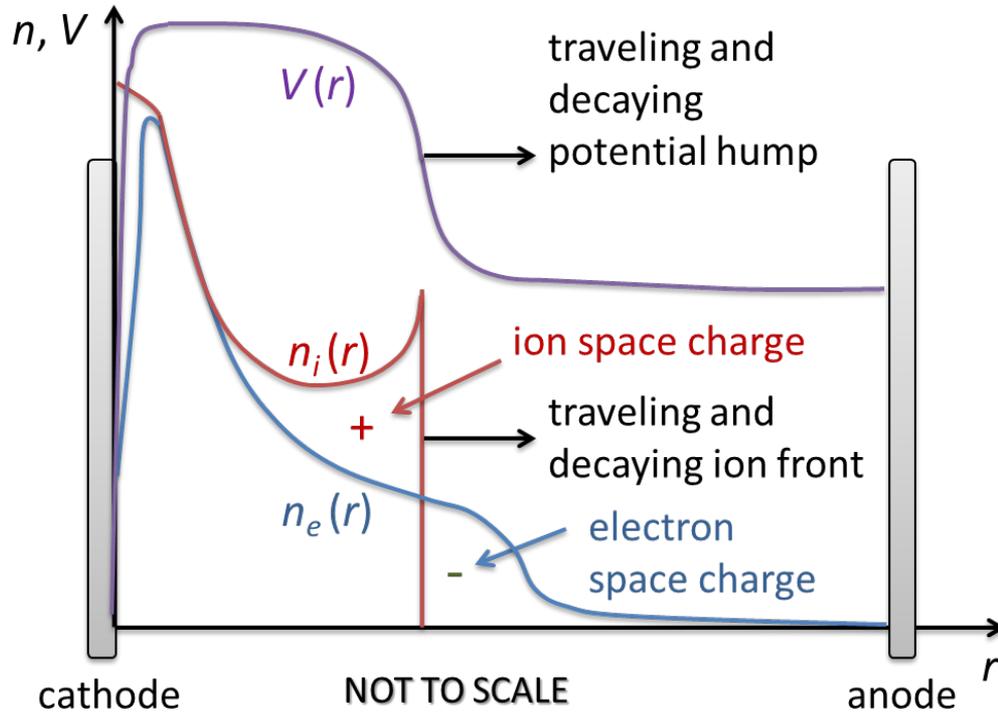
A. Anders and G. Yushkov, J. Appl. Phys. **91** (2002) 4824
E. Byon and A. Anders, J. Appl. Phys. **93** (2003) 1899

Ion acceleration near cathode spots: Potential hump vs. gas-dynamic theories



- Potential hump theory: ion acceleration in an electric field
- Gas-dynamic theory: acceleration by pressure gradients

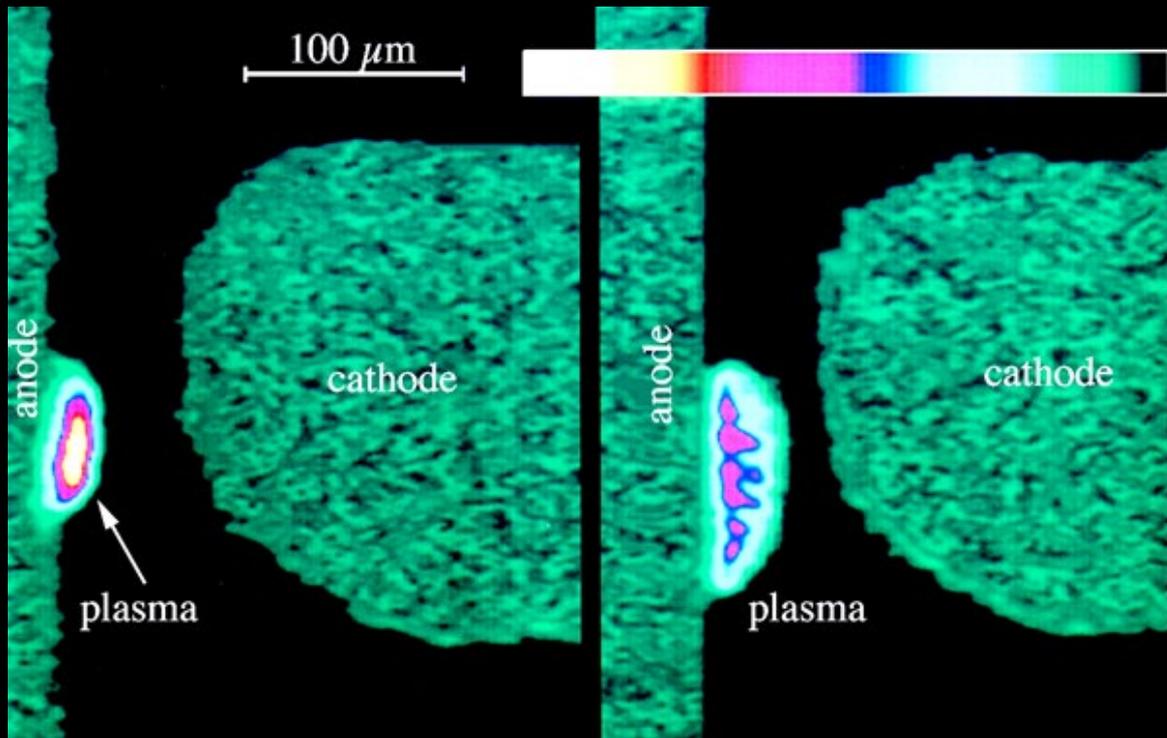
Both: Acceleration by gas-dynamic and electric field acceleration



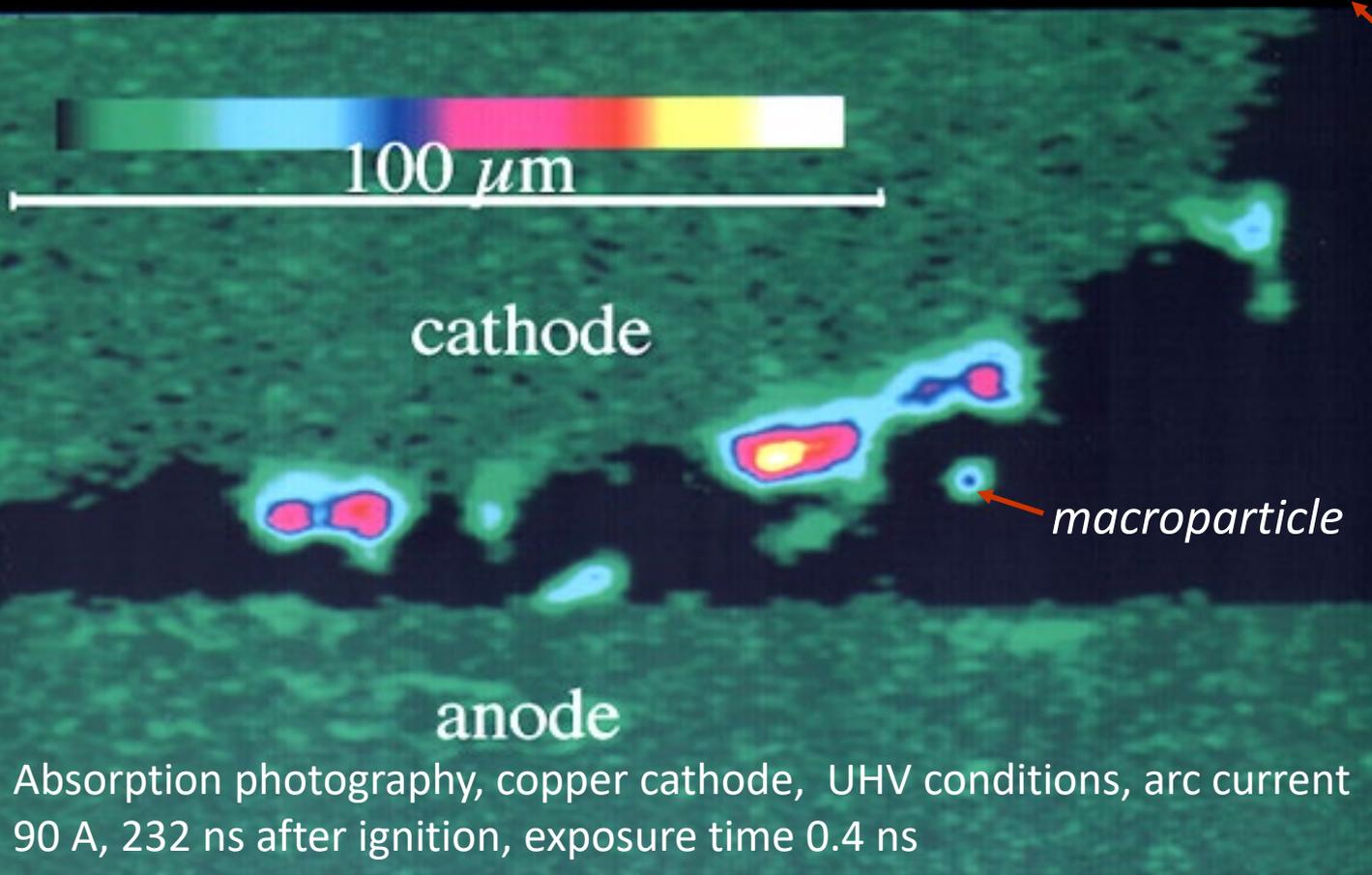
- Traveling and decaying potential humps
- Gradients are associated with the non-stationary nature of plasma formation
- Potential humps are stochastic
- Quantitative model remains to be developed

Cathodic Arc Plasmas: Observation of Gap Breakdown

- ❑ Laser absorption photography of vacuum breakdown event, Cu, 100 A, Δt between pictures 3 ns
- ❑ Spark phase of the arc: voltage is still high



Cathodic Arc Plasmas: Spot & “Spot Fragment” Formation



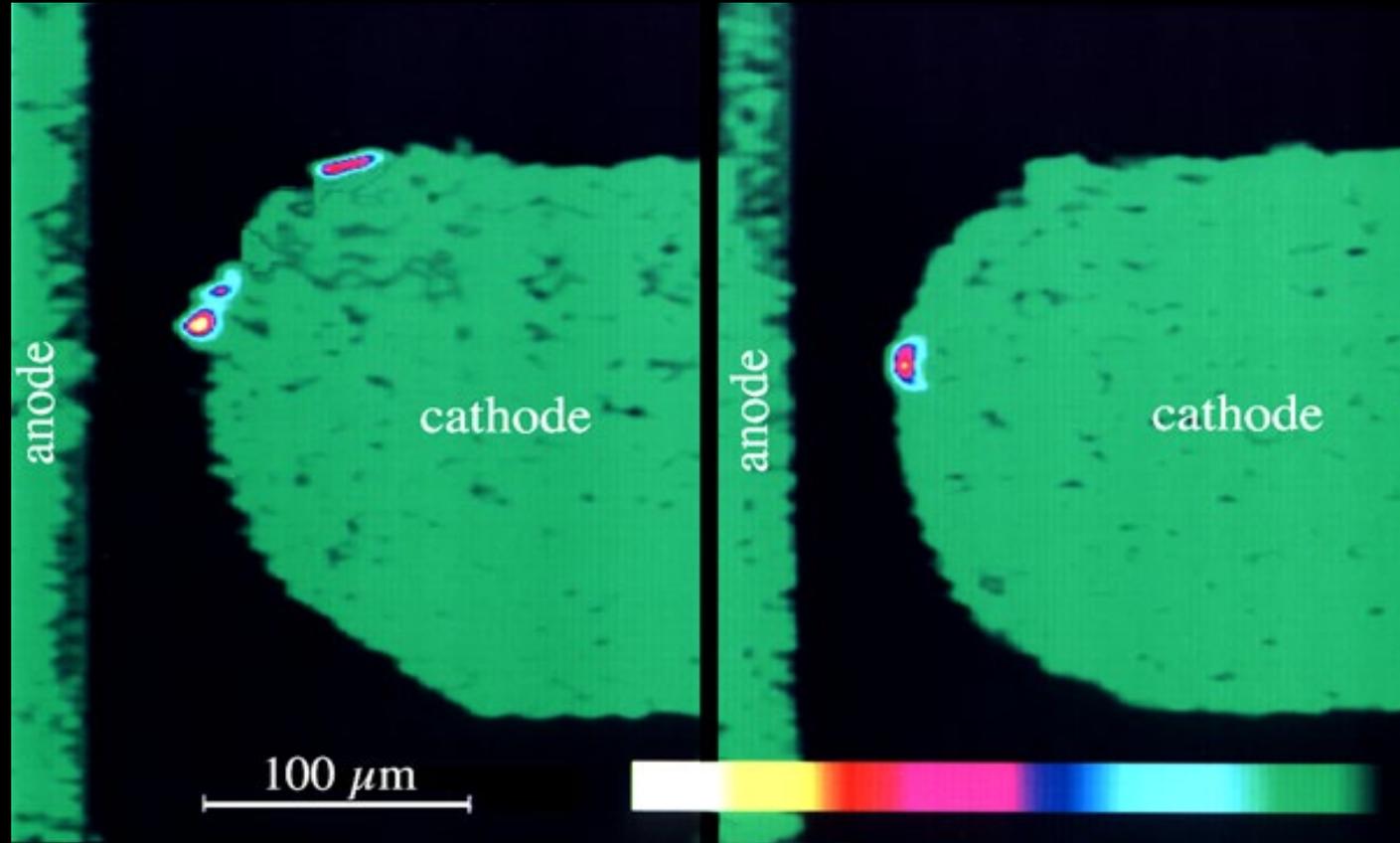
*Terminology
“fragment” not
needed if one
thinks of a fractal.*

*Example: visually
one spot but
containing
“fragments”*

Absorption photography, copper cathode, UHV conditions, arc current 90 A, 232 ns after ignition, exposure time 0.4 ns

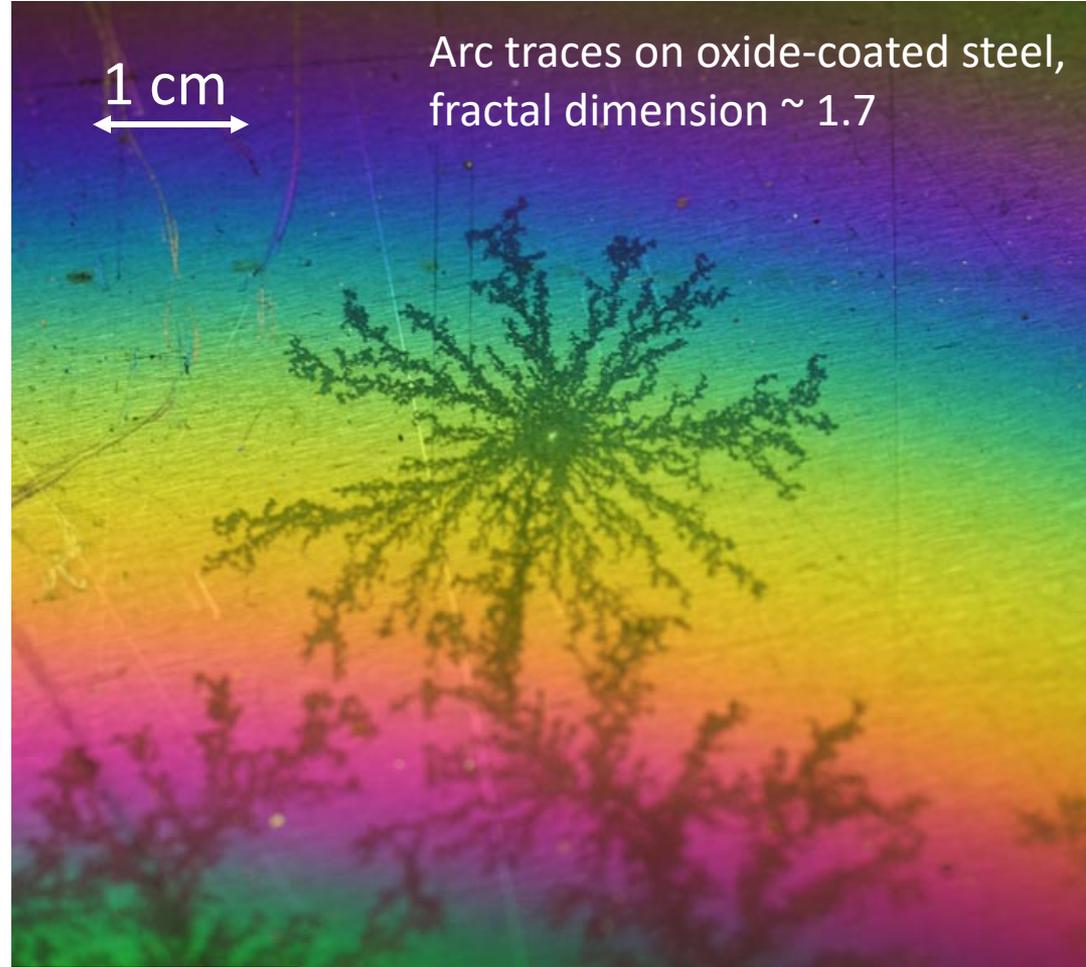
Cathode Spot Dynamics

development of
cathode spots,
observed by
absorption
photography, Cu,
100 A, Δt between
pictures 3 ns



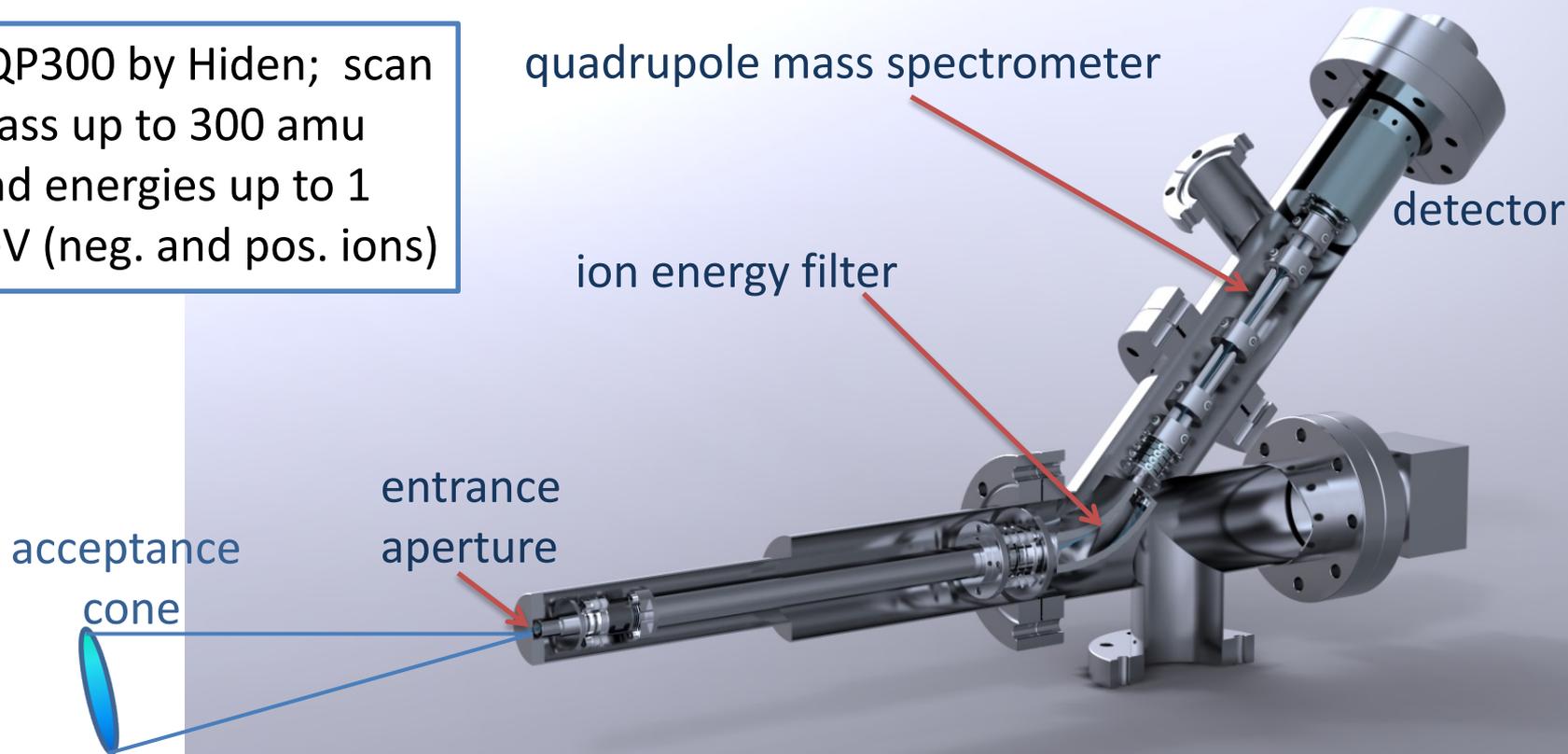
The Fractal Nature of Cathode Spots

- An object (“fractal”) is self-similar (= invariant with scaling) if it is reproduced by magnifying some portion of it.
- Self-similarity may be discrete or continuous, deterministic or probabilistic.
- Self-similarity can be mathematically exact or only approximate and asymptotical.
- Physical fractals have lower and upper scaling limits

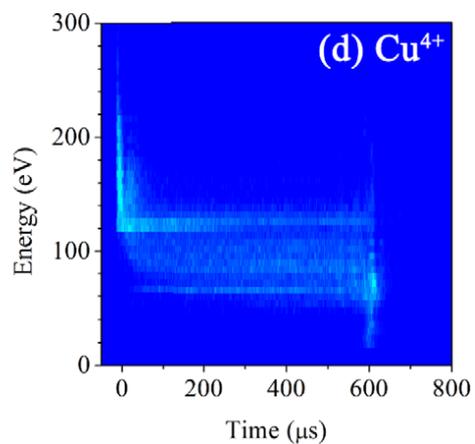
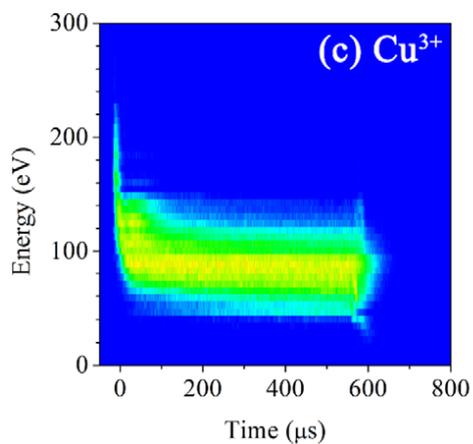
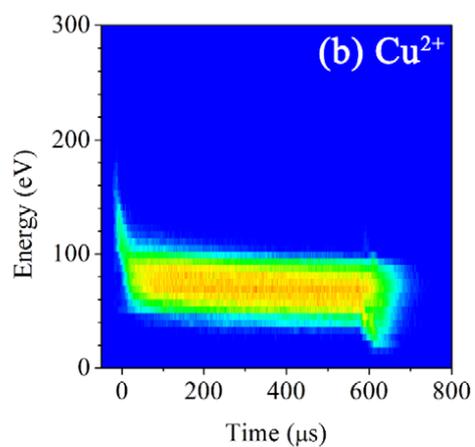
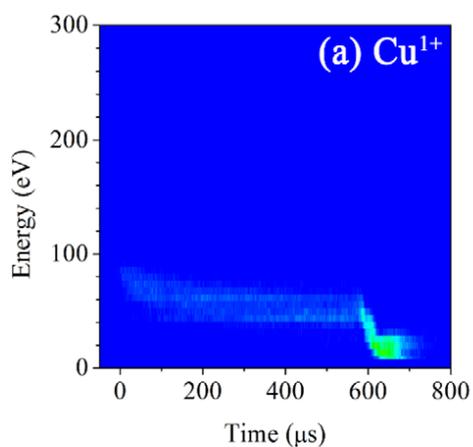


Ion flux spectrometry via ion energy/charge filter plus quadrupole mass/charge filter

EQP300 by Hiden; scan
mass up to 300 amu
and energies up to 1
keV (neg. and pos. ions)

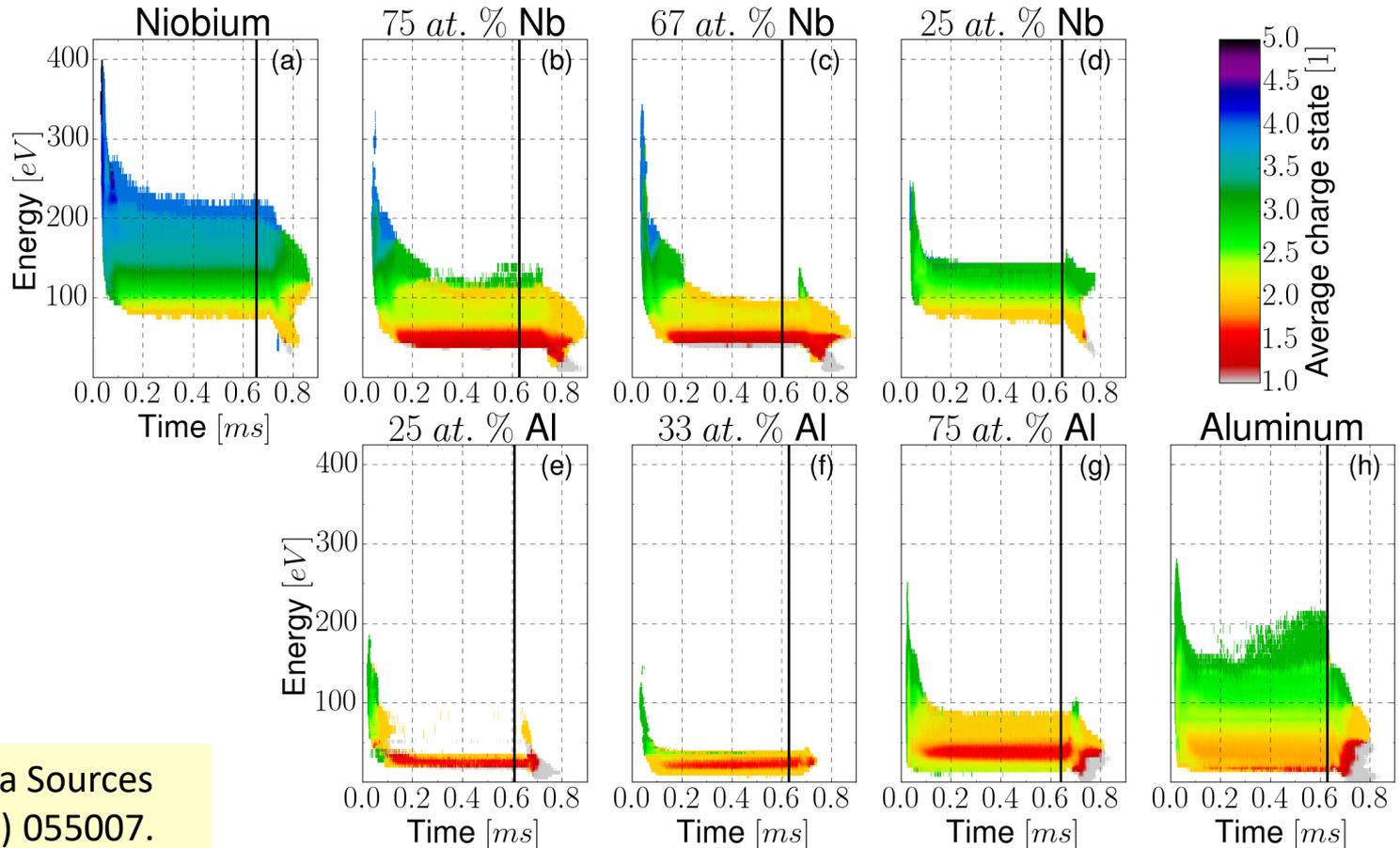


Charge-state and time-resolved energy distribution functions for cathodic vacuum arcs

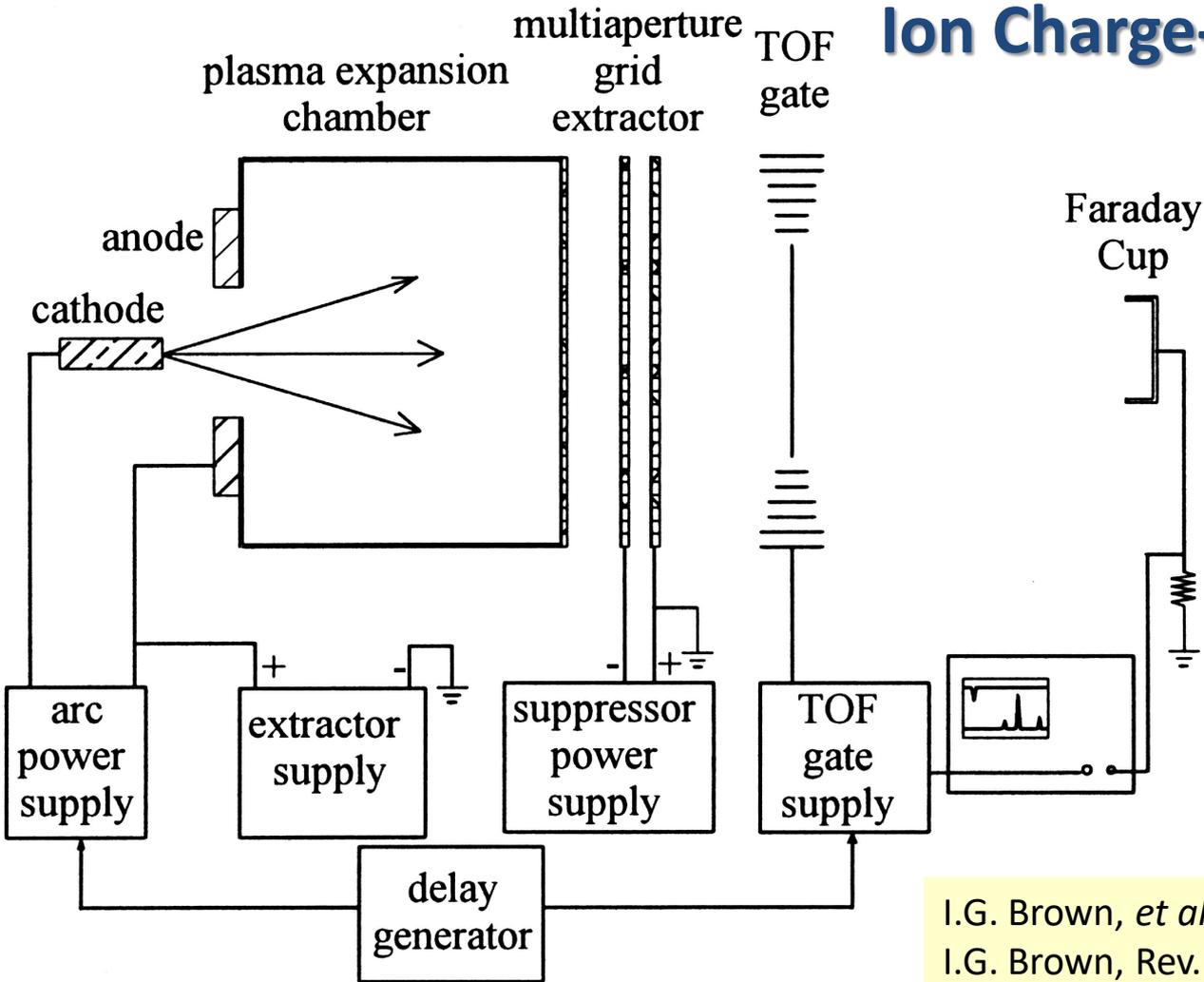


- multiply charged ions
- higher at beginning of discharge
- steady-state after $\sim 100 \mu\text{s}$
- fluctuating data, results show average over many measurements

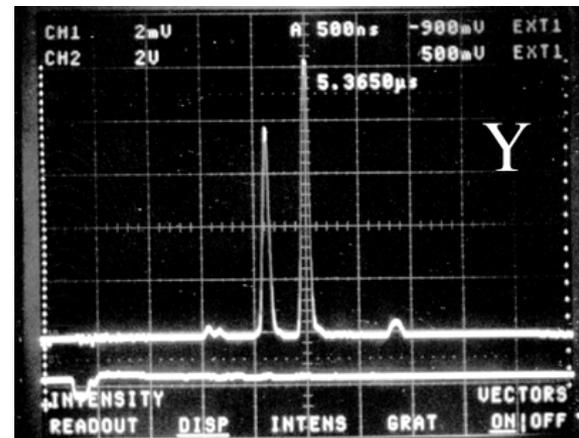
Charge-state and time-resolved energy distribution functions for cathodic vacuum arcs: Nb-Al cathodes



Ion Charge-State Spectrometry via TOF



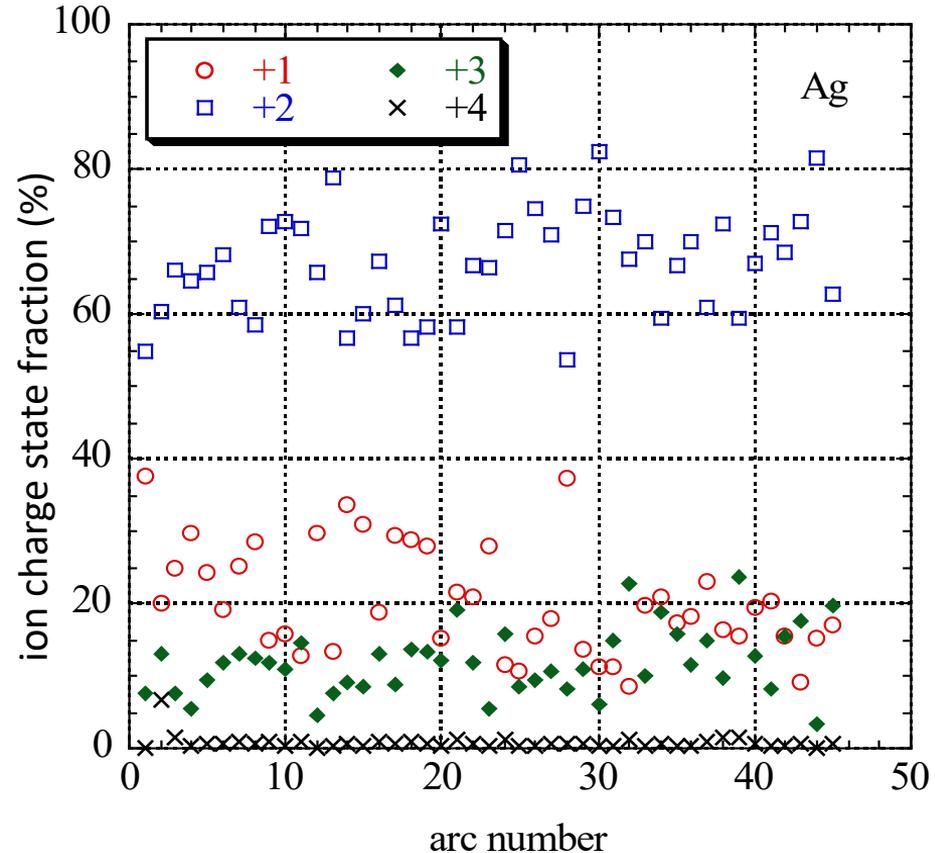
time-of-flight (TOF) method



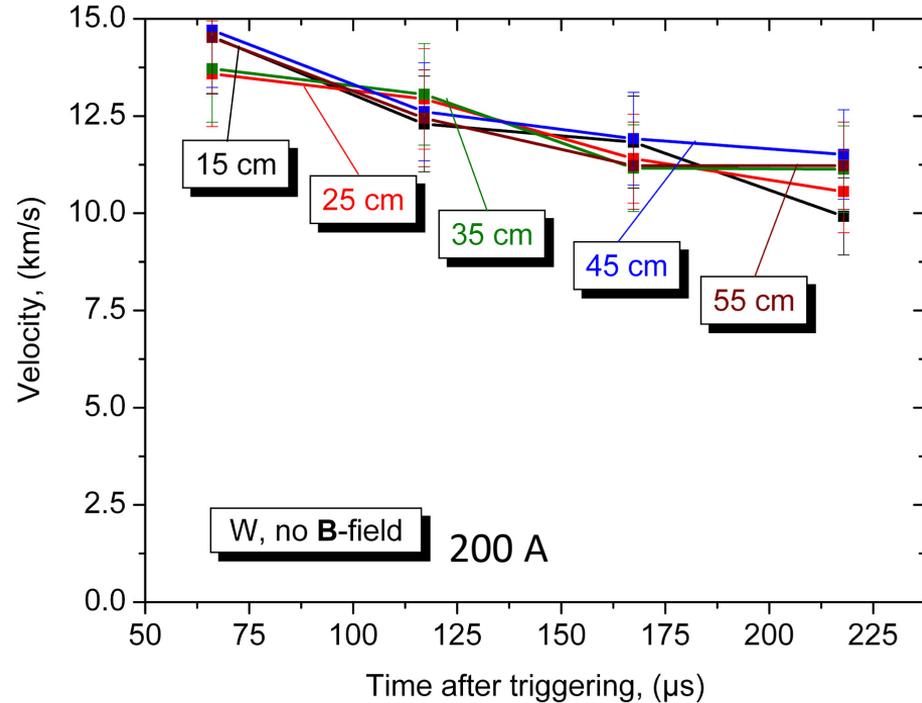
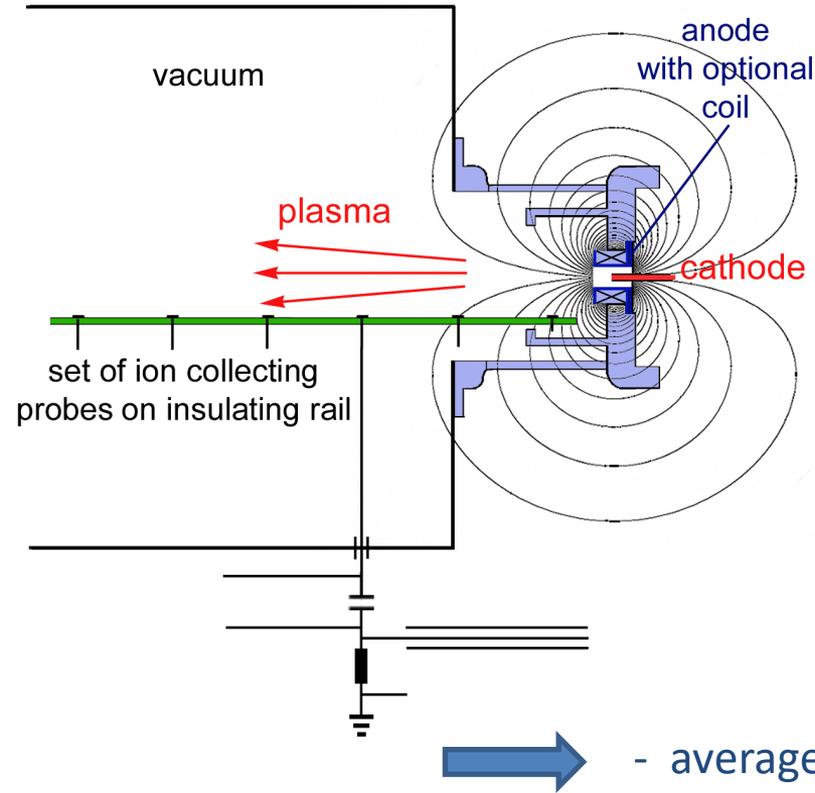
I.G. Brown, *et al.*, Rev. Sci. Instrum. 58 (1987) 1589.
I.G. Brown, Rev. Sci. Instrum., 65 (1994) 3061.

Fluctuations or “noise” can be observed in all parameters

- ❑ Fluctuations (or “noise”) is due to non-stationary plasma formation
- ❑ “Noise” is of little or no concern for plasma deposition (averaging effect) but essential for operation of the arc
- ❑ Example here: Ion charge state sampling of a pulsed silver arc



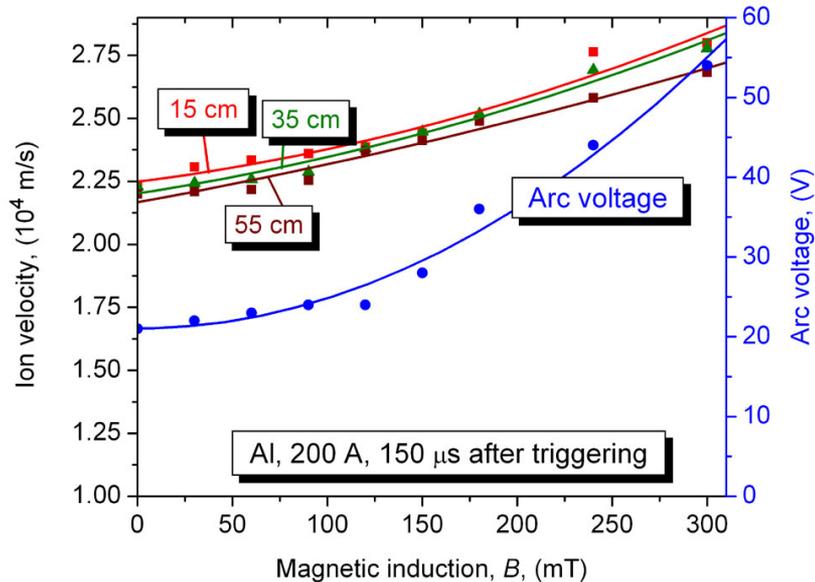
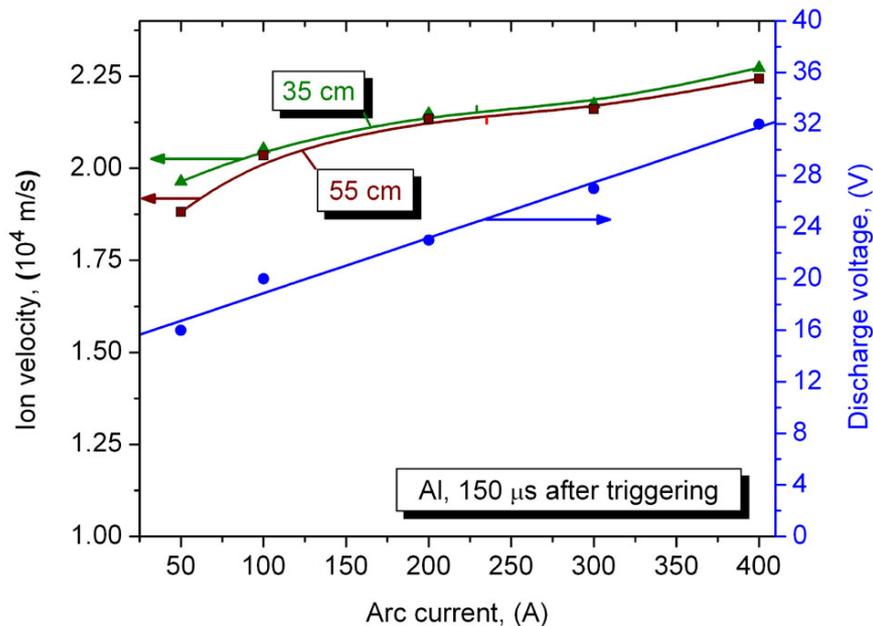
Evolution of ion velocities (time, distance from cathode)



- average velocity reduces slightly as time elapsed after arc start
- velocity is essentially constant as plasma propagates away from the cathode

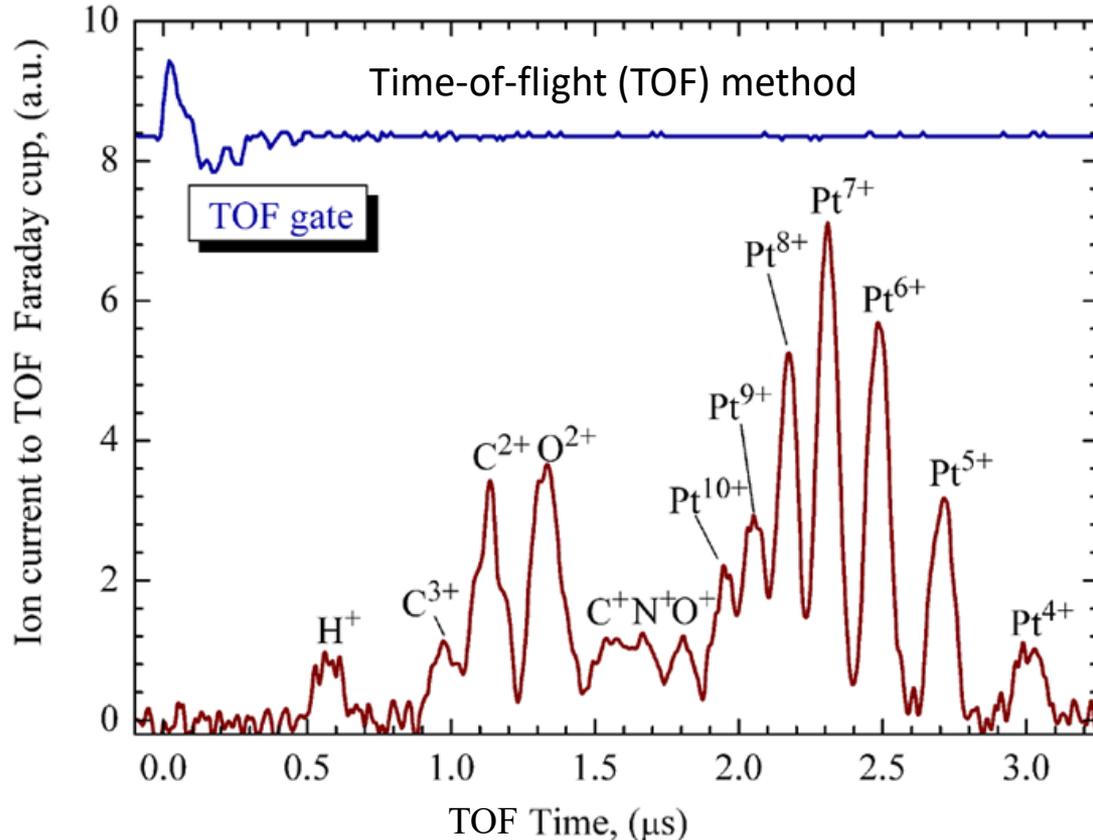
S. Hohenbild, et al., J. Phys. D: Appl. Phys. **41** (2008) 205210.

Evolution of ion velocities (arc current, magnetic field)



- ion velocity slight increases with arc current \rightarrow most of the acceleration happened at cathode spot, supplemented by acceleration in the near-cathode density gradients \rightarrow
- portion of additional near-cathode-acceleration can be enhanced by magnetic field (power enhancement)

Enhancing Ion Charge States



- 2+ and 3+ is typical for most metals
- using short pulses and a magnetic field \rightarrow up to 10^+ has been detected, indicative for high electron temperature and many collisions in dense plasma

G. Yushkov and A. Anders,
Appl. Phys. Lett. **92** (2008) 041502

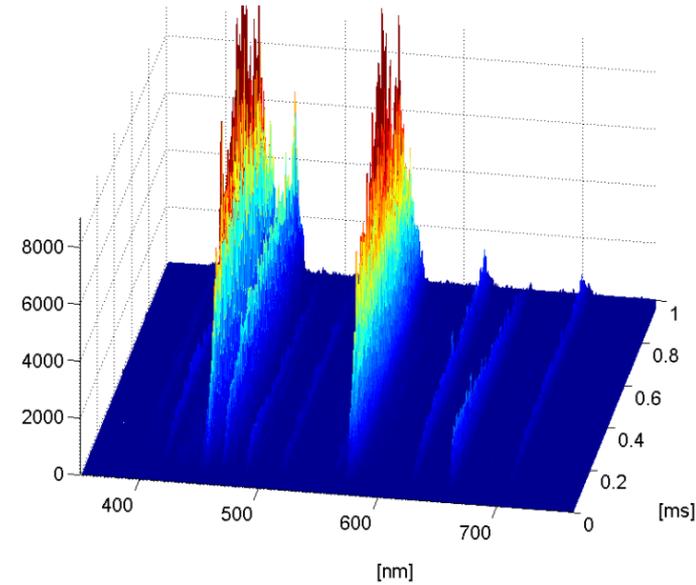
Sticking coefficient
is less than unity

Reducing Ion Charge States

neutrals lead to charge exchange
 $M^{Q+} + M^0 \rightarrow M^{(Q-1)+} + M^+$
and thus to a reduction of higher charge states

P.A. Ni, A. Anders, J. Phys. D: Appl. Phys. **43** (2010) 135201.

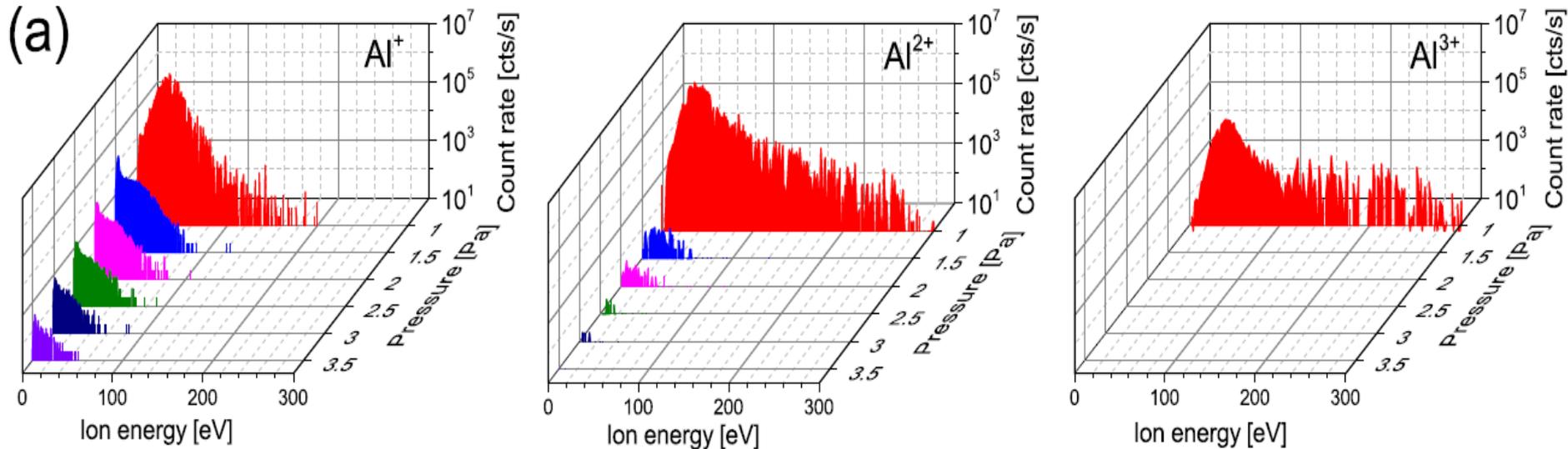
optical spectrum: Al neutral lines
appear and increase with time



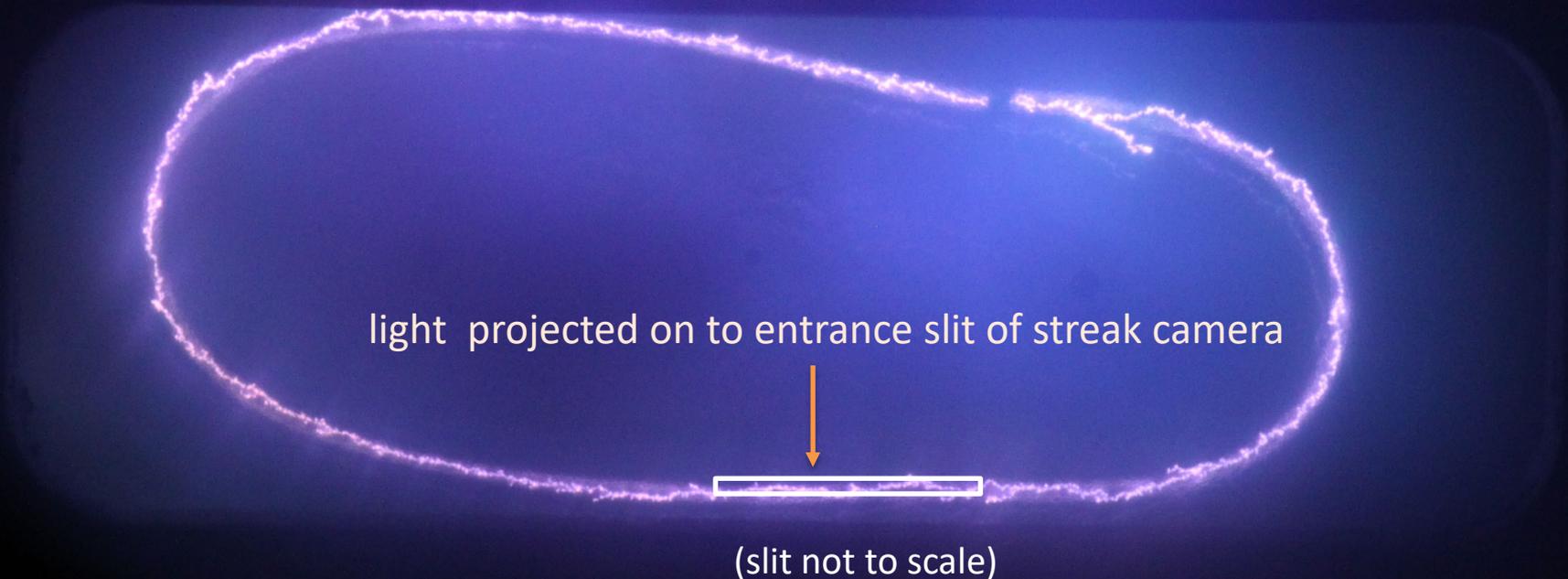
Introduction of reactive gas reduces ion charge states and ion energies

Example here: Al cathode in O₂ atmosphere, effects due to

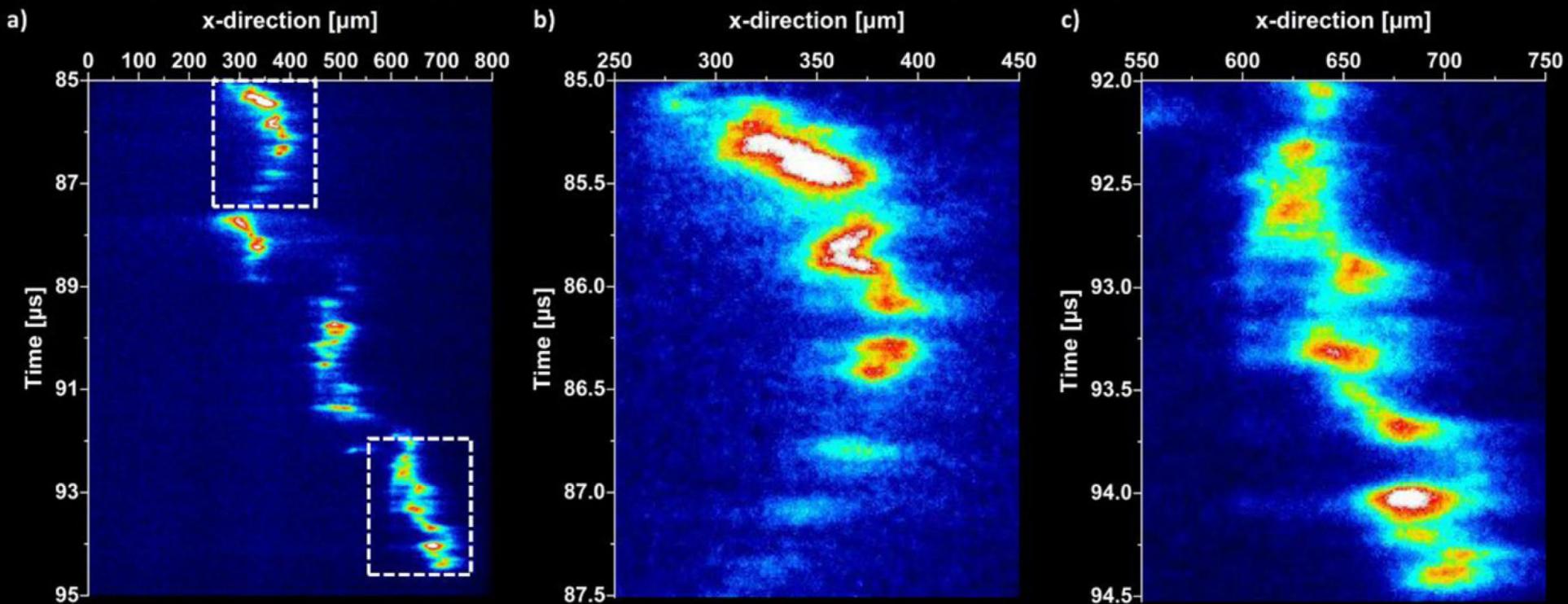
- collisions, mostly charge exchange collisions, and
- changes of the chemical composition of the cathode itself (“poisoning”)



Further studies of fluctuations using a “magnetically steered” arc: Search for “characteristic times



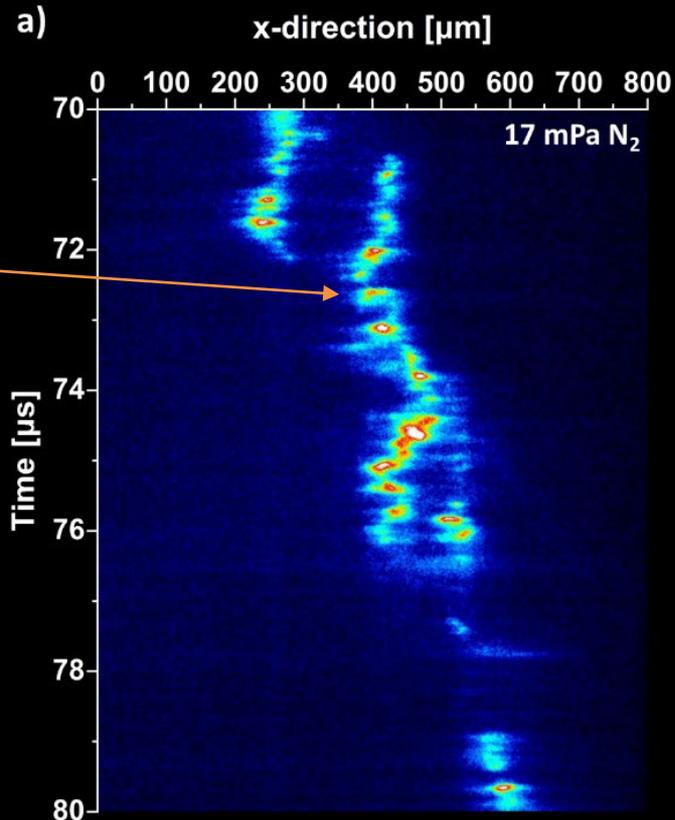
Study of Fluctuations: Cathode Spots are Fractals



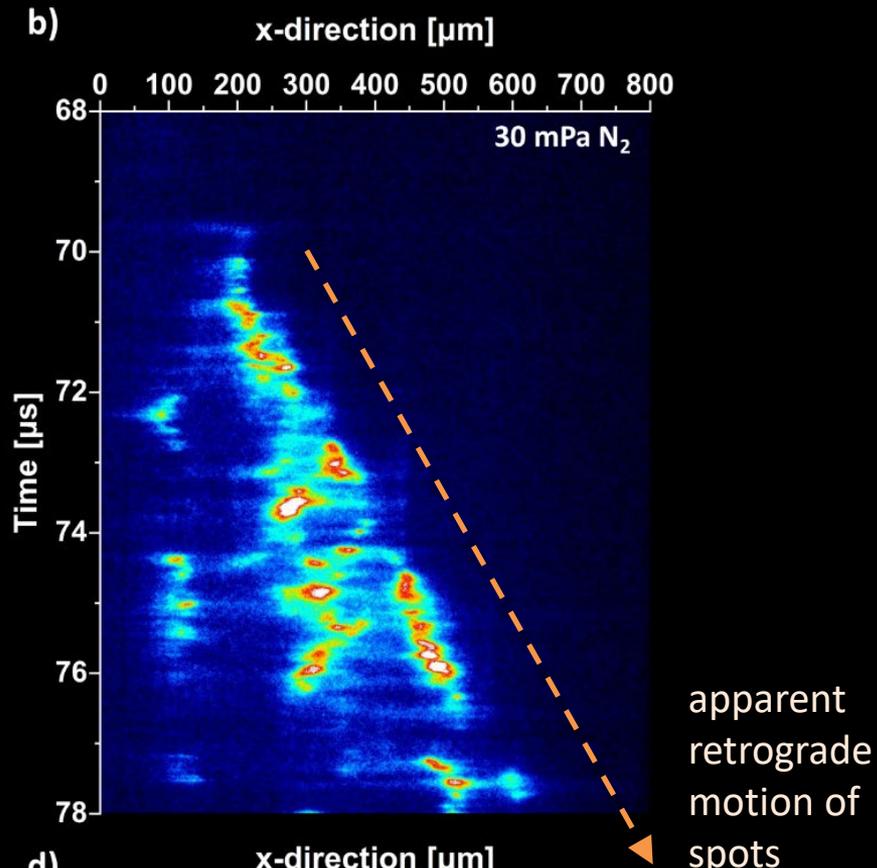
FFT analysis \rightarrow spots do not show periodic but self-similar (fractal) features

Study of Fluctuations: Cathode Spots as used in reactive arc-PVD

repeated
ignition at
about the
same site



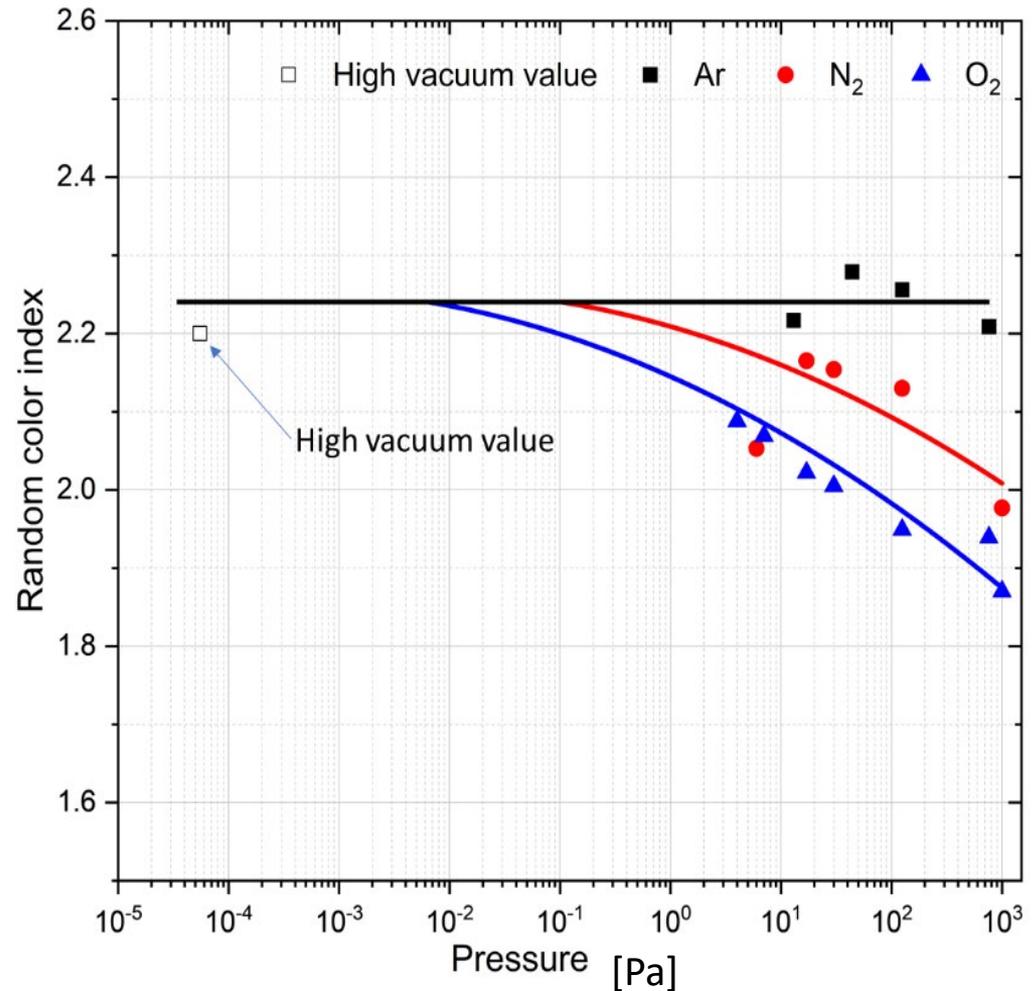
c) x-direction [μm]



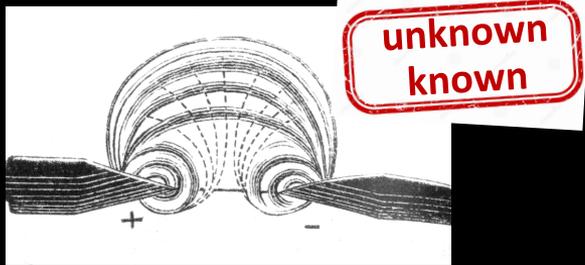
d) x-direction [μm]

Cathodic arc spots in the presence of a reactive gas

1. fluctuations or “noise” can be quantified by Random Color Noise (RCN) index
2. FFT shows no peaks \rightarrow no characteristic times of “arc cycles”
3. compound layers promote the ignition of new spots and reduce the feedback dependence on previous spots



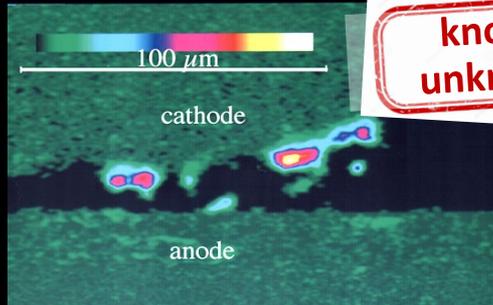
Summary



unknown
known

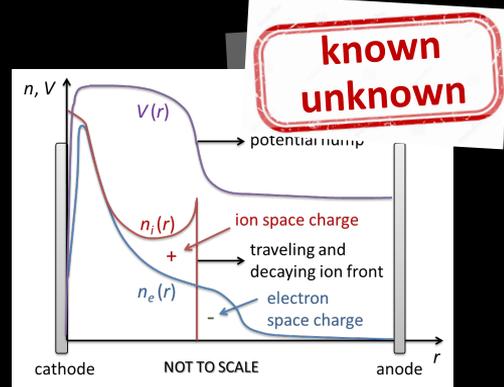
arcs and arc history

cathode spots



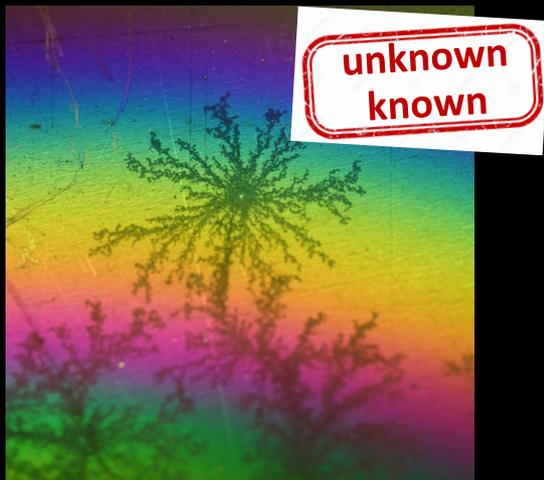
known
unknown

spot and "fragments"



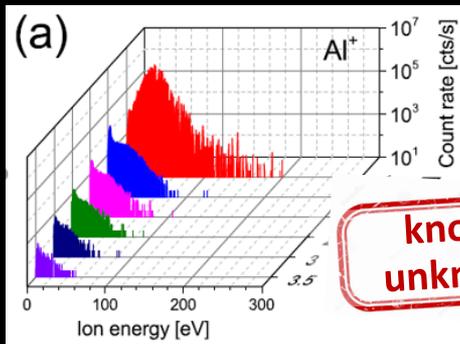
known
unknown

qualitative ion
acceleration model



unknown
known

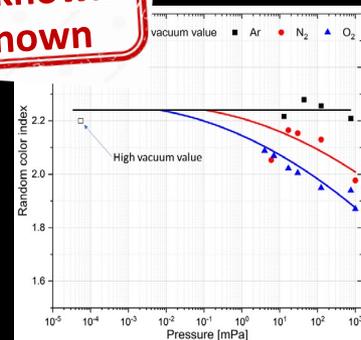
fractal nature of cathodic arcs



known
unknown

ion charge state and
energy distribution
functions

unknown
known



Quantification of fluctuations