The Interplay of Plasma Turbulence and Magnetic Reconnection in Producing Nonthermal Particles

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### Outline

- ► Is turbulence in magnetically-dominated (i.e., relativistic) plasmas an efficient source of nonthermal particles?
- ► If so, how does the nonthermal spectrum depend on the system parameters?
- Mechanism of particle acceleration? Interplay with magnetic reconnection?



# Nonrelativistic Turbulence (plus Reconnection)

▶ Plasma turbulence produces sheets of strong current density ⇒ natural sites of magnetic reconnection



(also Matthaeus & Lamkin 1986, Politano et al. 1989, Biskamp 2003, Mininni et al. 2006, Wan et al. 2013, Zhdankin et al. 2013)

# Nonrelativistic Turbulence (plus Reconnection)

Magnetic reconnection leads to a steepening of the energy spectrum. What about particle acceleration?



(also Carbone et al. 1990, Carbone 1995, Cerri & Califano 2017, Franci et al. 2017, Mallet et al. 2017, Loureiro & Boldyrev 2017, Walker et al. 2018, ..)

► Fully-kinetic treatment:

 $\Rightarrow$  we solve the coupled Vlasov-Maxwell system of equations through the PIC method

- ▶ PIC code TRISTAN-MP (Spitkovsky 2005)
- ▶ Decaying turbulence in relativistic pair plasmas
- ▶ 2D and 3D numerical simulations

# Tools and basic setup (Relativistic "Servidio Problem")



- mean magnetic field  $\langle \boldsymbol{B} \rangle = B_0 \hat{\boldsymbol{z}}$
- turbulence develops from uncorrelated magnetic fluctuations  $\delta B_x$  and  $\delta B_y$  in Fourier harmonics
- energy-carrying scale:  $l = 2\pi/k_f$ ( $k_f$  is the wavenumber where the energy spectrum peaks)
- Simulations in the magnetically-dominated regime

$$\sigma_0 = \frac{\delta B_{\rm rms0}^2}{4\pi w_0} \gg 1, \qquad \frac{1}{16} \le \frac{\delta B_{\rm rms0}^2}{B_0^2} \le 16, \qquad \theta = \frac{k_B T}{m_e c^2} \sim 1$$

with  $w_0 = nm_ec^2 + nk_BT \left[\hat{\gamma}/(\hat{\gamma} - 1)\right]$ 

## Fully-developed turbulence state

2D simulation with  $\sigma_0 = 10, \, \delta B_{\rm rms0}/B_0 = 1, \, L/d_{e0} = 3280$ 



▶ Copious presence of current sheets, plasmoids, and vortices

### Fully-developed turbulence state





Similar power-law particle energy distributions also in simulations by Zhdankin *et al.* 2017, 2018.





### Particle Spectrum: initial temperature scan



#### Particle Spectrum: system-size scan

2D simulations with  $\sigma_0 = 10, \, \delta B_{\rm rms0}/B_0 = 1$ 



### Current sheets in 3D



#### Current sheets in 3D: slices of $J_z$



## Current sheets in 3D: plasmoids/flux ropes



## Current sheets in 3D: plasmoids/flux ropes





#### How are particles accelerated?



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#### How are particles accelerated?



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▶ ~ 97% of the particles belonging to the nonthermail tail experience a sudden energy jump from the thermal pool



▶ Particle injection at reconnecting current sheets



▶ ~ 95% of the high-energy particles reside at injection at  $|J_{z,p}| \ge 2 J_{z,\text{rms}}$  (current sheets)



## Particle energization



► Typical particle energization history (2D)

### Particle energization



► Typical particle energization history (3D)

## Test particle spectrum



▶ For test particles that do not feel  $E_{\parallel}$ , there is a dramatic drop of particles in the nonthermal tail (while p is similar)

## Particle mixing





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## Energization mechanisms



▶ Particles energization occurs in a two-stage process

## Energization mechanisms: injection



### Energization mechanisms: stochastic acceleration





More details in Comisso & Sironi, PRL 121, 255101 (2018)

#### Particle Spectrum: particles per cell scan

2D simulation with  $\sigma_0 = 10, \, \delta B_{\rm rms0}/B_0 = 1, \, L/d_{e0} = 820$ 



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# Plasmoid formation: different spatial resolutions



From two 2D simulations with 3 cells per  $d_{e0}$  (top row) and 10 cells per  $d_{e0}$  (bottom row).

## Movie - time evolution of $J_z$

