

Microturbulence-mediated route for stronger energetic ion transport and Alfvénic mode intermittency in ITER-like tokamaks

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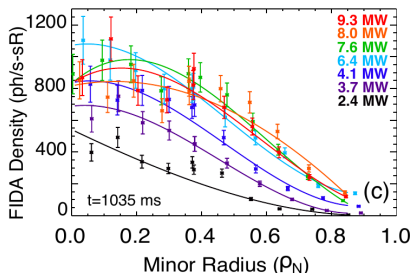
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Classical α -particle confinement & Coulomb collisions imply weak Alfvénic mode driven α -transport in burning plasmas

α s, beam confinement is crucial for self-sustained plasma with current drive and auxiliary heating.

- Recent Critical Gradient DIII-D XPs imply stiff EP transport with resilient EP profiles (Collins et al., PRL'16) \rightarrow
- M. Fitzgerald et al., NF'16 (see also M. Schneller et al., PPCF'16) predicted $n=1-35$ unstable Alfvénic eigenmodes (AEs) at $\delta B_r/B \lesssim 3 \times 10^{-4}$ driven by fusion alphas in ITER. HAGIS was used to assess α -particle relaxation in the presence AEs.
- Amplitudes are too low to produce a noticeable α -transport in ITER (5% of all alphas: Fasoli, TPB, NF'07).
- 1 MeV NBI ion drive was ignored.
- AE amplitude increase by a factor ~ 50 is required to see appreciable alpha losses.



We are showing that the anomalous pitch angle scattering such as due to the micro-turbulence is required for higher AE amplitudes.

Resonance Broadened Quasi-linear (RBQ) model computes time dependent EP dynamics

Action-angle formalism using flux variables results in a set of equations for fast ion DF:

(Kaufman,PhFI'72, Berk,Breizman, NF'95, adapted for RBQ Duarte, PhD'17, Gorelenkov, NF'18, APS'18)
 RBQ quasilinear equations include 3 time scales: γ_L^{-1} , γ_d^{-1} , and v_{eff}^{-1}

$$\frac{\partial}{\partial t} f = \pi \sum_{l,k} \frac{\partial}{\partial P_\phi} C_k^2 \mathcal{E}^2 \frac{G_{m'l}^* G_{mp}}{|\partial \Omega_l / \partial P_\phi|_{res}} \mathcal{F}_l[\Omega] \frac{\partial}{\partial P_\phi} f + v_{eff}^3 \left| \frac{\partial \Omega_l}{\partial P_\phi} \right|^{-2} \frac{\partial^2}{\partial P_\phi^2} (f - f_0),$$

where distribution f is evolved due to scattering terms on the RHS and AE growth rate.

AE amplitudes evolve due to explicit equation

$$C_k(t) \sim e^{(\gamma_L + \gamma_d)t} \Rightarrow \frac{dC_k^2}{dt} = 2[\gamma_L(t) + \gamma_d] C_k^2.$$

*AE growth rates γ_L are evolved due to $f(t)$; γ_d are fixed.

Critical for RBQ multi-mode cases (Dupree'66, Berk'95, White'18) is

the resonant frequency and its broadening (nonlinear bounce ω_{bNL} and effective scattering v_{eff}):

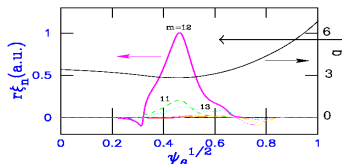
(V. Duarte APS'20 inv. talk)

$$\delta(\Omega = \omega + n\dot{\phi} - m\dot{\theta} - l\omega_b) \rightarrow \text{window function, } \mathcal{F}_l[\Omega]$$

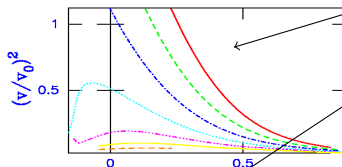
RBQ provides EP diffusion to TRANSP whole device modeling code to evolve EP distribution function.

Extensive V&V are undertaken within ISEP SciDAC. Self-consistent diffusion near the resonance region (Duarte AAPPs, APS'20 on Friday).

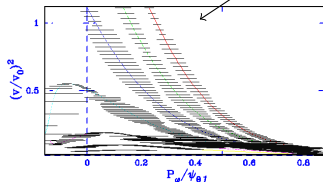
RBQ workflow for $n = 4$ Reversed Shear Alfvén Eigenmode (near q_{min})



- Ideal MHD code NOVA finds RSAE structure $f = 84\text{kHz}$ (Collins, PRL '16).
 - This mode provides a channel for ion diffusion and hollow fast ion pressure profiles: resonant particles are close to the injected pitch angle.



- NOVA-K code computes resonances for particle interactions with the mode and $\langle \mathbf{v} \cdot \mathbf{E} \rangle$ matrices.
- RBQ broadens those resonances along P_ψ direction using QL prescriptions for each mode.



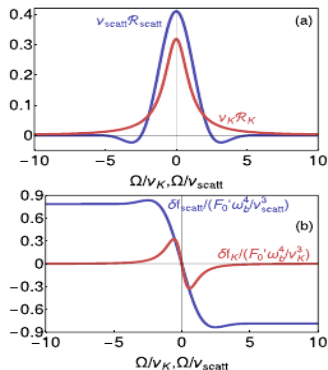
- Quasi-linear methodology is adopted for its numerical efficiency.
 - Rigorous V&V is implied through ISEP SciDAC.

Shown is the broadening at measured amplitude $\delta B_\theta / B = 7 \times 10^{-3}$.

Selfconsistent wind. function accounts for radial (or P_ϕ) diffusion (Duarte PoP'19)

A new QL theory is developed for near threshold regimes basing on near threshold WF:

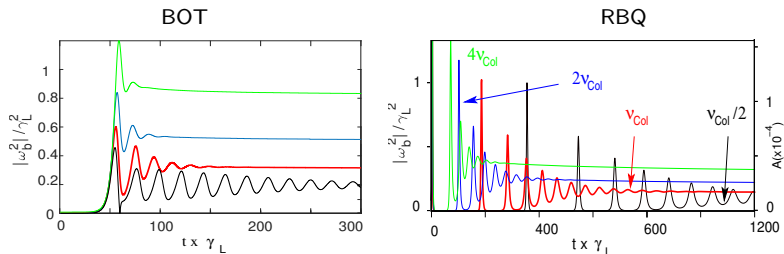
$$0.3 \lesssim v_{\text{eff}} / (\gamma_L - \gamma_d) < 1, \quad \gamma_d / \gamma_L > 0.35.$$



- QL methodology accounts for resonant dynamics *selfconsistently*.
- Measurable interplay of growth/damping rate scales is important for experiments.
- Energy slowing down is weak and is included in recent work: Duarte, Lestz et al., in preparation (see J. Lestz poster, Wed. PP12).
- Window function is required for ITER for efficient and (yet) realistic simulations to account for α -particle dynamics near the resonances with AEs.

RBQ+near threshold window function & Bump-on-tail (BOT) agree qualitatively

RBQ (Quasi-linear): fixed boundary conditions (BC) at the axis, fixed (zero) at the edge.
 BOT (Vlasov kinetic equation solver): Ω , γ_L , γ_d , v_{eff} , fixed BC at infinity.



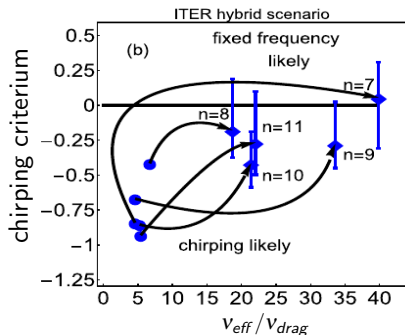
Nominal case for $n = 3$, $f = 75\text{kHz}$: $\gamma_d = -0.75\gamma_L = -2.075\%$ at $v_{Coul} = 8.9\text{sec}^{-1}$ or $v_{eff} = 8.017 \times 10^3\text{sec}^{-1}$.

In RBQ the recovery time (repetition rate) scales for nominal DIIID case (*Collins et al., PRL'16*) as $\Delta t (\simeq v_{eff}^{-1}) = 19\text{msec}$. In BOT it is $\sim 30\text{msec}$.

(*Gorelenkov & Duarte, PLA'20 in press*)

Expectations for ITER plasmas are based on earlier chirping studies

Coulomb scattering can be enhanced by the microturbulence in ITER



V.Duarte, N.Gorelenkov et al., NF'18

- Projections of microturbulence from $E_\alpha = T_i$ to $E_\alpha = 3.52 \text{ MeV}$ are made with the help of Zhang et al., PRL'08.
- Pitch angle scattering is increased by the factor 4 to 10. Expected AE amplitude can grow to $A \sim v_{eff}^2 = (16 \div 100) v_{Coul}^2$.
- Fitzgerald et al., NF'16 cite ~ 50 times required AE amplitude increase to see appreciable α losses.
- No NBIs were included.

- Gyrokinetic simulations are needed to compute robustly pitch angle scattering in ITER experimental conditions to make reliable predictions.

Summary

- 1 AE instabilities can lead to losses and significant redistribution of alphas and beam ions in ITER if mediated by microturbulence.
 - 1 Coulomb scattering alone is not strong enough for strong losses.
 - 2 Microturbulence needs to be accounted for.
- 2 At the moment 2D version of RBQ is being developed within ISEP SciDAC for realistic simulations:
 - 1 RBQ is to be used for burning plasma applications, ITER, where multiple instabilities are expected.