

Zero-frequency zonal flows and predatorprey behavior driven by Alfvén modes in the JET tokamak

J. Ruiz Ruiz¹, J. Garcia², M. Barnes¹, M. Dreval³, V. Duarte⁹, C. Giroud⁴, V.H. Hall-Chen⁵, M. Hardman⁶, J.C. Hillesheim⁷, Y. Kazakov⁸, S. Mazzi², F.I. Parra⁹, B. Patel⁵, A. Schekochihin¹, Z. Stancar⁵
 1. Oxford 2. CEA 3. KIPT 4. UKAEA 5. A*STAR 6. Tokamak Energy 7. CFS 8. LPP-ERM/KMS 9. PPPL

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Why care about **energetic particles** (EPs) and **turbulence** in magnetic confinement fusion?



• Energetic particles (α 's) are ubiquitous to fusion energy: D + T $\rightarrow \alpha + n$

 $E_{\alpha} = 3.5 \text{ MeV} >> T_{\text{plasma}} \sim 20 \text{ keV}$

- Today's experiments use external EP sources for heating (NBI, ICRH, E~100 keV)[*]
- When the plasma is self-heated by α 's, we call it a **burning plasma**



Today's experiments



Burning: No further need of external heat!

- To produce energy from fusion, need to satisfy Lawson criterion (1955): $n_p \tau_E \gtrsim const.$
- Confinement time τ_E is determined by **turbulence**!

[*] Neutral beam injection / Ion cyclotron resonance heating

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Today's experiments are predominantly heated with ~ 100 keV particles. How are MeV energetic particles different from 100 keV (1/2)?





- If $v_{\rm EP} > v_A \rightarrow$ EPs can resonate with Alfvén waves as they slow down
- MeV EPs more easily destabilize Alfvén waves (Alfvén eigenmodes AEs)
 Greatly decreases the confinement of the energetic particles

How are MeV energetic particles different from 100 keV (2/2)?

- 3. Have larger orbits \rightarrow more difficult to confine Vertical cross-section of a tokamak orbit excursion δr in JET (B = 3T, a \approx 90 cm): $\delta r(3.5 \text{ MeV}) \sim 10 - 20 \text{ cm}$ $\delta r(100 \text{ keV}) \sim 2 - 4 \text{ cm}$ а 4. Rotation (NBI~100 keV): shears away turbulence [*] but α -particle source is isotropic \rightarrow no rotation $v \neq 0$ (NBI) $v = 0 (\alpha' s)$ Sheared turb. Unsheared turb. \rightarrow Smaller eddies,

[*] Terry RMP 2000

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- Bad news for MeV energetic particles in burning plasmas?
- 1. Critical energy \rightarrow e- heating, but need ions hot for fusion!
- 2. $v_{\alpha 0} > v_A \rightarrow$ easier to destabilize Alfvén modes : worse confinement
- 3. Larger orbits \rightarrow worse confinement of EPs
- 4. Isotropic α -particle source \rightarrow no rotation expected in burning plasmas (ITER, SPARC, STEP, DEMO)





...



5

EPs can have a positive effect on confinement by affecting the dominant linear micro-instability (ITG) and nonlinearly, the turbulence

- 1. Dilution:
 - Only background ions contribute to the ion-temperature-gradient driven mode (ITG)
 - Quasineutrality: less resonant ions \rightarrow stabilize ITG [1] $Z_i n_i + Z_{EP} n_{EP} = n_e$
 - Usually not a dominant effect $(n_{EP} \ll n_i)$
- 2. Shafranov's shift (β') stabilization:
 - Energetic particles contribute to β'
 - → can affect MHD equilibrium & micro-instabilities [2]
- 3. More recently, active linear kinetic effect of energetic particles can stabilize ITG [3]

1,2,3 are *linear*, and extrapolate to be small in a burning plasma

 $\beta_s' = \sum \left(\frac{a}{L_{Ts}} + \frac{a}{L_{ns}} \right) \beta_s$

4. Nonlinear effects: "Nonlinear electromagnetic stabilization" [4], presence of EPdriven modes [5], zonal flows ...

[1] Tardini NF 2007 [2] Bourdelle PoP 2003 [3] Di Siena NF 2018/PoP 2019, Wilkie NF 2018 [4] Citrin PRL 2013 [5] Di Siena NF 2019, Mazzi Nat Phys 2022, Biancalani JPP 2023

Experiments in JET with MeV ions have exhibited enhanced thermal confinement that could not be explained by linear theories [1]

- Confined MeV ions (ICRH) destabilized Toroidal Alfvén eigenmodes (TAEs). Observed improved confinement [2]
- Nonlinear gyrokinetic (turbulence) simulations showed that TAEs generate zonal flows that stabilize the turbulence
 Decrease in the thermal ion heat transport





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These experiments show that MeV energetic particles can have an unexpected, beneficial effect on the thermal plasma 😉

[1] Tardini NF 2007, Bourdelle NF 2005, DiSiena NF 2018/Wilkie NF 2018 [2] Using 3-ion ICRH scheme: Mazzi Nat. Phys. 2022, Ruiz Ruiz PRL 2025 Juan Ruiz Ruiz | Zero-frequency zonal flows and predator-prey behavior in JET The JET experiments probe close-to-burning plasmas, and exhibit nonlinear interplay between energetic particles and the thermal plasma

- Achieved by the presence of confined MeV ions
 - Provided e- heating
 - Destabilized Toroidal Alfvén eigenmodes (TAEs)
 → improved ion confinement through zonal flow
- Generation of zonal flow by TAEs is a nonlinear process predicted theoretically and numerically [1]

TAE + TAE \rightarrow ZF(freq~0)TAE + TAE \rightarrow sidebands(freq ~ 2x, 3x, ..)



 Stationary, zero-frequency zonal perturbation could interact with and stabilize turbulence → consistent with Mazzi Nat Phys 2022 / Ruiz Ruiz PRL 2025

• Zero-frequency fluctuation driven by AEs lacks experimental confirmation We show the first experimental confirmation of a zero-frequency fluctuation that is pumped by an Alfvén eigenmode in a magnetically-confined plasma³

[1] Todo NF 2010 (MHD) / Chen PRL 2012, Qiu PoP 2016/NF 2017 [2] Biancalani EPS 2019 [3] Ruiz Ruiz PRL 2025 Juan Ruiz Ruiz | Zero-frequency zonal flows and predator-prey behavior in JET The JET experiments probe close-to-burning plasmas, and exhibit nonlinear interplay between energetic particles and the thermal plasma

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 The effect of α's can be deleterious () or beneficial (): will the overall effect

 of alpha particles on the thermal plasma be beneficial for fusion performance?

 [1] Todo NF 2010 (MHD) / Chen PRL 2012, Qiu PoP 2016/NF 2017

 [2] Biancalani EPS 2019

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Can the effect of MeV ions be beneficial to global confinement?



- Experiments probing close-to-burning plasmas in JET
 - JET plasma with MeV-range ICRH-heated ions
 - Doppler backscattering (DBS) measurements
 - Power spectrum and bicoherence analysis (Ruiz Ruiz PRL 2025)
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JET 97090 L-mode heated with ICRH produces MeV range 3He that slows down on electrons

• L mode,
$$I_p$$
=2.4 MA, B_0 =3.2 T, $\bar{n}_{e,l} = 7 \times 10^{19} \text{m}^{-2}$

- ICRH heating (no NBI) steps:
 - 2 MW (low), 4 MW (medium) and 7 MW (high)
- H+D (background), ~0.2% trace 3He → 4-5 MeV [*]
- Fast 3He slows down on e- (>90% heating)
- ➔ Almost pure electron heating plasma via slowing down of 3He mimic conditions of 3.5 MeV alpha particles in a burning plasma



[*] '3-ion ICRH heating', Kazakov NF 2015

[*] Nocente NF 2020, Kazakov PoP 2021, Mazzi Nat Phys 2022

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[**] Beurskens NF 2021

Confinement improves and ion temperature increases with ICRH power

- Core electron temperature T_e increases as expected
- L-mode confinement factor H_{89,P} increases
 (consistent with [*]) → confinement improvement

- Ion temperature ALSO increases in deep core -- puzzle
 - e- and ions more collisionally decoupled as T_e increases
 - No *T_i* clamping [**]
 - Ion turbulence is stabilized?





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 - → Decrease in Q_i ($Q_e > Q_i$, TRANSP)



High P_{ICRH}

15

16

17

14

t [s]



Med P_{ICRH}

13

12

Te.axis [keV]

Low P_{ICRH}

11

[n_d/ [1020m]

Increase in T_i and thermal confinement is correlated with increased Alfvénic activity observed with increasing ICRH power





- Low ICRH power similar to Ohmic
- Medium ICRH power close to marginal stability of MHD modes
- High ICRH power MHD activity (TAE)
 f ≈ *f*_{TAE} ≈ 270 300 kHz
 n ≈ 4 6

$$f_{\rm TAE} = v_A / 4\pi q R$$

- Electron heating (MeV-range fast ions)
- Improved ion confinement (T_i , $H_{89,P}$ increase)
- Alfvénic activity

Motivate local measurements of turbulence fluctuations

Perform turbulence measurements using Doppler Backscattering (DBS) diagnostic





- Externally launched microwave beam propagates into plasma until it encounters a cutoff.
- Forward beam deviated upwards.
- **Detect** backscattered radiation from turbulence wavenumber **k**
- Scattered power $P_s \propto \langle |\delta n(\mathbf{k}_1)|^2 \rangle$
- DBS measures one k_{\perp} at one location (cutoff) [*].



[*] Hall-Chen PPCF 2022

In what follows: Show local turbulence measurements at $\rho \approx 0.3$ from $k_{\perp}\rho_s \approx 0.5 - 4$

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DBS measurements at medium ICRH power





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At medium P_{ICRH} , DBS beam reaches deep core, spectrogram exhibits periodic bursts and spectral peaks at gap frequency f_{TAE} and harmonics





- $f_{\text{TAE}} \approx v_A / 4\pi q R \approx 270 300 \text{ kHz}$
- Spectral peaks at f_{TAE} , $2f_{TAE}$, $3f_{TAE}$

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Low P_{ICRH} Broadband fluctuations (~ITG) **Medium P_{ICRH}** Broadband background + spectral peaks At medium P_{ICRH} , DBS beam reaches deep core, spectrogram exhibits periodic bursts and spectral peaks at gap frequency f_{TAE} and harmonics Spectrogram $P_s(f,t)$



- $f_{\text{TAE}} \approx v_A / 4\pi q R \approx 270 300 \text{ kHz}$
- Spectral peaks at f_{TAE} , $2f_{TAE}$, $3f_{TAE}$
- Presence of TAEs (near q=1, tornado TAEs [*]) that burst every $\Delta t \approx 6 7$ ms

[*] Saigusa PPCF 1998

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- $f_{\text{TAE}} \approx v_A / 4\pi q R \approx 270 -$ 300 kHz
- Spectral peaks at $f_{\text{TAE}}, 2f_{\text{TAE}}, 3f_{\text{TAE}}$

 $^{-1}$

-2

- Presence of TAEs (near q=1, tornado TAEs [*]) that burst every $\Delta t \approx 6 -$ 7 ms
- Are higher gap modes $2f_{\text{TAE}} \bullet$ sidebands?
 - Faint effect on • turbulence (ITG)?

[*] Saigusa PPCF 1998

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At medium P_{ICRH} , bicoherence exhibits three-wave matching relations at f_{TAE} , $2f_{TAE}$, $3f_{TAE}$ and a zero-frequency fluctuation



$$|B_{s}(f_{1},f_{2})| = \frac{\langle \hat{A}_{j}(f_{1})\hat{A}_{j}(f_{2})\hat{A}_{j}(f_{1}+f_{2})^{*} \rangle}{\langle |\hat{A}_{j}(f_{1})\hat{A}_{j}(f_{2})|^{2} \rangle^{\frac{1}{2}} \langle |\hat{A}_{j}(f_{1}+f_{2})|^{2} \rangle^{\frac{1}{2}}}$$

→ three-wave phase matching relationship between f_1 and f_2 such that $f_1 \pm f_2 = f_3$ [*]

Mode-mode interactions (sidebands!):

$$(f_1, f_2) = \mathbf{A}: (f_{TAE}, f_{TAE}) \rightarrow 2f_{TAE}$$
$$\mathbf{B}: (2f_{TAE}, -f_{TAE}) \rightarrow f_{TAE}$$

• Interaction with a zero-frequency fluctuation: $(f_1, f_2) =$ **C**: $(2f_{TAE}, 0) \rightarrow 2f_{TAE}$ **D**: $(2f_{TAE}, -2f_{TAE}) \rightarrow 0$ At medium P_{ICRH} , bicoherence exhibits three-wave matching relations at f_{TAE} , $2f_{TAE}$, $3f_{TAE}$ and a zero-frequency fluctuation

0.64

0.56

0.48

0.40

0.32

0.24



$$|B_{s}(f_{1},f_{2})| = \frac{\langle \hat{A}_{j}(f_{1})\hat{A}_{j}(f_{2})\hat{A}_{j}(f_{1}+f_{2})^{*} \rangle}{\langle |\hat{A}_{j}(f_{1})\hat{A}_{j}(f_{2})|^{2} \rangle^{\frac{1}{2}} \langle |\hat{A}_{j}(f_{1}+f_{2})|^{2} \rangle^{\frac{1}{2}}}$$

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• Mode-mode interactions (sidebands!): $(f_1, f_2) = \mathbf{A}: (f_{TAE}, f_{TAE}) \rightarrow 2f_{TAE}$

To our knowledge: first experimental measurement of three-wave phase matching relationship between Alfvénic modes and a zero-frequency fluctuation in a magnetically confined plasma

D: $(2f_{\text{TAE}}, -2f_{\text{TAE}}) \rightarrow 0$

DBS measurements at high ICRH power





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 Modes in higher gaps (up to 5f_{TAE}) inside q=1 surface (and persist outside)





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- Bursts of the Alfvénic modes (Δt ≈ 4 ms). Are higher gap modes sidebands?





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Causality?





- Modes in higher gaps (up to 5f_{TAE}) inside q=1 surface (and persist outside)
- Bursts of the Alfvénic modes $(\Delta t \approx 4 \text{ ms})$. Are higher gap modes sidebands?
- Power in turbulence frequencies (f < 200 kHz) decreases during bursts (predator-prey behavior?)

TAE in fundamental gap f_{TAE} precedes modes in higher gaps and turbulence stabilization \rightarrow suggests that fundamental TAE is the drive of the sidebands, affecting the turbulence (zonal flow?)



The DBS bicoherence shows that the strongest phase-matching

• Strongest interactions involve $(2f_{TAE}, 0)$ and $(3f_{TAE}, 0)$, not $(f_{TAE}, 0)$



The JET experiments suggest that MeV ions can be *beneficial* to the global confinement (via Alfvén eigenmodes and the zero-frequency fluctuation)

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Linear CGYRO at high P_{ICRH} shows destabilization of fast-ion-driven instability consistent with AE, and ITG





- Highly driven TAE ($k_y \rho_s < 0.1$), ITG ($k_y \rho_s [0.3 1]$)
 - CGYRO $\omega_{\text{TAE}} \approx 4 c_s/a$ matches DBS $f_{\text{TAE}} \approx 270 300 \text{ kHz}$, $n \approx 4 6 (k_y \rho_{sD} \approx 0.04)$
- Negligible effect of EPs on ITG \rightarrow *linear* effects (dilution, β' , kinetic stabilization) are negligible
- No low- k_y AE without EP

Low k_y mode (AE) exhibits critical gradient behavior with $a/L_{T_{\text{fast}}}$





- $k_y \rho_{sD} \approx 0.04 \text{ mode (AE) matches:}$
 - DBS freq. ($f_{\text{TAE}} \approx 260 290 \text{ kHz}$)
 - $n \approx 4 6 (k_y \rho_{sD} \approx 0.04)$
 - $\delta E_{||} = 0$
- $a/L_{T_{\text{fast}}} < 8$: long-tail mode (~MT)
- $a/L_{T_{fast}} > 8: AE$
- → Probe turbulence without fast particles, at low, medium, and at high a/L<sub>T_{fast} drive of AE.
 </sub>

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Local nonlinear CGYRO for high P_{ICRH}: turbulence is stabilized when AE is driven unstable





- Electromagnetic $(\delta \phi, \delta A_{||})$, fast-Maxwellian 3He $(T_f \approx 168 T_i)$
- AE+ITG scales: $k_y \rho_{sD} = [0.02, 1.26], L_y = 314 \rho_{sD}$ $k_x \rho_{sD} = [0.015, 2.93], L_x = 410 \rho_{sD}$
- Heat flux dominated by ions for no EPs, $Q_i/Q_{gB} \approx 3$
- Stabilization of $Q_{i,e}$ with unstable AE $a/L_{T_{\text{fast}}} \gtrsim 4$ (marginal AE instability)

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Density power spectrum and bicoherence at $a/L_{Tfast} = 6$ (low drive) exhibits harmonics of f_{TAE} and dominant mode-mode interactions



Density power spectrum and bicoherence at $a/L_{Tfast} = 8$ (med drive) exhibits harmonics of f_{TAE} and dominant interaction with f = 0



Conclusions from the research (1/3)



To our knowledge: the first experimental confirmation of a zero-frequency fluctuation that is pumped by an Alfvén eigenmode in a magnetically-confined plasma [*]. Confirmed by local nonlinear gyrokinetic simulations.

- JET discharge with dominant e- heating (MeV range fast ions) shows T_i and $H_{89,P}$ increase with P_{ICRH} , decrease in $Q_i \rightarrow$ enhanced global confinement.
- Alfvénic modes (tornado TAEs) exhibit nonlinear phase-matching relations with a zero-frequency fluctuation, and suggest predator-prey behavior TAE-turbulencezonal flow.
- Nonlinear CGYRO confirms nonlinear phase-matching between f=0 and AE mode → zero-frequency fluctuation could be responsible for enhanced total plasma confinement, could balance deleterious energetic particle transport by the Alfvén eigenmodes.

[*] Ruiz Ruiz et al., PRL 2025

Discussion of the findings (2/3)



- What this work means
 - Zero-frequency zonal flow pumped by Alfvén modes is real [1]. We see a manifestation of it in bicoherence and the predator-prey behavior.
 - Possibly a glimpse of the rich, nonlinear behavior at the heart of next-generation burning plasmas.
- Why has this not been seen before?
 - MeV ions: difficult to generate
 - Most machines do not confine MeV ions (except JET, TFTR, JT-60U)
 - NBI can drive Alfvén modes, but also often stabilize them [2]
 - Here, we have analyzed a discharge exclusively heated by MeV ions
 - Zonal flow from AEs never really been looked after until very recently

[1] Todo NF 2010/2012, Qiu PoP 2016/NF 2017, Mazzi Nat Phys 2022

Important / unresolved questions (3/3)



- 1. Measure radial structure of zero-frequency zonal flow produced by Alfvén modes: need further measurements!
- 2. What does this zonal flow depend on? Can it be tailored to our advantage in a future burning plasma? eg. increase Ti?
- 3. Are Alfvén eigenmodes good or bad overall for fusion performance?
- 4. Here, MeV ions are controlled externally what will happen in a burning plasma, where nonlinear effects are expected to be more important?